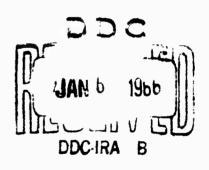
GD/C-DBE-65-028 Space Science Laboratory

# EMPIRICAL INFRARED ABSORPTION COEFFICIENTS

of  $\rm H_2O$  from 300°K to 3000°K

C. C. Ferriso
C. B. Ludwig
and
A. L. Thomson

DECEMBER 1965



Technical Report

Research and Advanced Technology Department

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#### **ABSTRACT**

A set of absorption coefficients for water vapor as functions of temperature has been empirically dedecued from existing quantitative absorption and emission spectra between 1 and 22 \( \mu\$. The basic assumption was made that, for optically thick gases, the curve of growth is given by a statistical model. The absorption coefficients were obtained for each vibration-rotation band by adjusting the band averaged line half widths and line spacings such that the integral of the absorption coefficients was equal to the known integrated intensity of the particular band at a given temperature. The effect of foreign gas broadening and non-resonant water molecule broadening was included. The present results are compared with independent theoretical and experimental results and the agreement is satisfactory.

#### INTRODUCTION

For the purpose of analyzing and predicting rocket plume radiation, it is necessary to have available a representation of the spectral properties of all the major radiating species as functions of the relevant parameters. This representation should be as simple as possible yet detailed enough to allow for determining the essential features of the spectrum.

The major infrared molecular radiators in present rocket exhausts are carbon dioxide, water vapor, and several diatomic molecules. The spectral properties of carbon dioxide and many of the diatomic molecules have been reported in the literature. However, there are at present no theoretical predictions of the complete water vapor spectra in the infrared region at high temperatures. It is the purpose of this paper to derive a set of absorption coefficients for water vapor between 50 and 7500 cm<sup>-1</sup> in the temperature range from 300° to 3000°K from available experimental data.

During the last few years, a large number of spectra of high temperature water vapor between 1 and 22 microns have been published. 1-12 These spectra have been measured by different techniques and cover the temperature range from 300° to 2700°K. The optical depths range between about 0.2 to 100 cm·atm, total pressures range between 50 mmHg and 10 atm, but the range of parameters covered at any one temperature is rather limited, especially at higher temperatures. These measurements show that the spectral emissivity depends on the following independent parameters: absorption coefficient, wavelength, temperature, partial pressure, total pressure, line broadening ability of foreign gases, and pathlength. The

most important parameter is the absorption coefficient as a function of wavelength and temperature. If sufficient thin-gas spectra at all temperatures were available, a set of absorption coefficients could be derived directly. However, this is not the case. In order to obtain absorption coefficients from measured spectra in which the gas was not thin, the functional relationship between emissivity or absorptivity and optical depth (curve of growth) must be known. A number of theoretical models for the curves of growth have been developed. 13,14 Goody 14 has treated the absorptivity of a gas whose spectral lines are distributed statistically both in position and intensity. In the same paper the applicability of the statistical model to water vapor was established. Also, Howard, Burch, and Williams 15,16 have applied this model to water at room temperature and 1 atm pressure and found good agreement over a wide range of optical depths.

The expression for the spectral curve of growth contains the mean spectral absorption coefficient, the optical depth, and a rotational fine structure term which is dependent on a collision parameter and the mean line spacing.

We propose to use these curves of growth to determine the mean spectral absorption coefficients from non-thin water spectra by adjusting the fine structure term in such a way that the integral of the absorption coefficients over a given vibration-rotation band results in the known value of the band intensity.

#### DETERMINATION OF ABSORPTION COEFFICIENTS FROM NON-THIN GAS SPECIRA

In the statistical model the spectral emissivity is given by

$$\varepsilon(\omega) = 1 - \exp\left[-\frac{W(\omega)}{d(\omega)}\right] \tag{1}$$

where W(w) is the average equivalent width and d(w) is the average line spacing. Goody has shown that the curves of growth are quite insensitive to the choice of the line intensity distribution functions. When the smoothed absorption coefficient and mean line width to line spacing ratio are chosen so that the curves of growth agree in the optically thin and in the square root limits, the values of  $\frac{W}{d}$  corresponding to a delta function, an exponential or a 1/S distribution differ by at most 25%.

In the present reduction of the experimental data, we have used the intermediate form corresponding to an exponential distribution of line intensities:

$$\frac{W(\omega)}{d(\omega)} = \frac{k(\omega)u}{\sqrt{1 + \frac{k(\omega)u}{4a(\omega)}}}$$
(2)

where k(w) is the mean absorption coefficient, u the optical depth (at STP), and a(w) the fine structure parameter. The fine structure term a(w) is proportional to the local mean value of the ratio of the collision half-width to the line spacing,  $\gamma/d$ . It is to be noted that the value of d, deduced from experimental curves of growth, is only meaningful when interpreted in terms of a particular intensity distribution function. In the present analysis, we have chosen to normalize the value of d so that it is appropriate for a delta function distribution (i.e.,  $a = \gamma/d$ ).

The line width is expected to be proportional to the number of collisions experienced by the water molecule per unit time. Since

collisions with other molecules of water vapor are more effective than those with molecules of a different species, it has been found that good agreement with the experimental data at room temperature has been obtained by assuming that the line width is proportional to  $(\gamma_{H_2} o^p_{H_2} o^p_{$ 

$$a(w) d(w) = p_{H_2O}(\gamma_{H_2O}^{O}(w) \frac{T_O}{T} + \gamma_{H_2O}^{O}(w) \sqrt{\frac{T_O}{T}}) + p_x \gamma_x^{O}(w) \sqrt{\frac{T_O}{T}},$$
 (3)

where  $\gamma_{H_2O}'$  is the line width due to non-resonating dipole interactions of  $H_2O$  at  $T_0$ . The following assumptions are now made: The line widths  $\gamma^O(\omega)$  due to interactions with resonating and non-resonating water molecules and with foreign gas species, which are functions of frequency, can be approximated by a band averaged value, so that the frequency dependence of the broadening term  $a(\omega)$  appears only through the mean line spacing  $d(\omega)$ . Thus, Eq. (3) becomes

$$a(\mathbf{w}) = \frac{\overline{\gamma_{H_2O}^o}}{\overline{d(\mathbf{w})}} p_{\overline{T}} \sqrt{\frac{T_o}{T}} \left[ c \left( \sqrt{\frac{T_o}{T}} + \overline{\sigma}' \right) + (1-c)\overline{\sigma}_{x} \right]$$
 (4)

where  $\overline{\sigma}' = \overline{\gamma_{H_2}'^{\circ}} \sqrt{\gamma_{H_2}^{\circ}} 0$  and  $\overline{\sigma}_x = \overline{\gamma_x'^{\circ}}/\gamma_{H_2}^{\circ} 0$  and c is the mole fraction of water. Values of  $\overline{\gamma_{H_2}^{\circ}} 0$  and  $\overline{\sigma}_x$  are given in the literature and the following values will be used in the present study:  $\overline{\gamma_{H_2}^{\circ}} 0 = 0.5 \text{ cm}^{-1} \text{atm}^{-1}$ ;  $\overline{\sigma}_{N_2} = 0.18$ ;

 $\sigma_{0}^{2} = 0.09$ . Since no value for  $\sigma'$  is given, we will assume initially that  $\sigma'$  is approximately 0.1. As a next step we assume that a band averaged mean line spacing d can be used so that a(w) becomes independent of the frequency and we can write Eq. (2) as

$$\frac{\mathbf{W}}{\mathbf{d}}(\mathbf{w}) = \frac{\mathbf{k}(\mathbf{w})\mathbf{u}}{\sqrt{1 + \frac{\mathbf{k}(\mathbf{w})\mathbf{u}}{\mathbf{k}_{\mathbf{n}}}}}$$
 (5)

The frequency-independent broadening term a can then be evaluated from existing spectra, where  $\frac{W}{d}(w) = -\ln (1-\epsilon[w])$ , by using as the second of the two equations required to define k(w) and  $\overline{a}$ 

$$\int k(\mathbf{w})d\mathbf{w} = \alpha^{0}\phi(\mathbf{T}), \qquad (6)$$

where  $\alpha^{\rm O}$  is the known integrated intensity at STP of the band in question and  $\phi({\bf T})$  is a known function of the temperature. <sup>21</sup> It can be shown that, in the temperature range of interest here (up to ~ 3000°K), the function  $\phi({\bf T}) \equiv 1$  for fundamental bands in general. However, as pointed out recently by W. S. Benedict, <sup>22</sup> the 6.3- $\mu$  fundamental band of  ${\bf H}_2$ 0 may not obey this rule because of strong rotation-vibration interaction. Room temperature measurements of the 6.3- $\mu$  band indicate a value of 250 (cm<sup>-1</sup> per cm at STP), while measurements at temperatures of 2700°K indicate a value of 380 (cm<sup>-1</sup> per cm at STP). Therefore, an empirical function  $\phi({\bf T}) = 250 + 0.05{\bf T}$  has been used to describe the temperature dependence of the 6.3- $\mu$  band between 300°K and 3000°K. For combination and overtone bands,  $\phi({\bf T})$  was given in Ref. 21 for the 1.87- and 1.38- $\mu$  bands. The values of  $\alpha^{\rm O}$  in (cm<sup>-1</sup> per cm at STP) for the four major vibration-rotation bands used in the present analysis are shown in Table I.

Table I: Integrated intensities used in present analyses.

Bands	6.3 µ	2.7 <u>µ</u>	1.87 µ	1.38 µ
$\alpha_{\mathbf{i}}^{\mathtt{c}}$	250 ± 20%	230 ± 15%	26 ± 15%	21 ± 15 <b>%</b>

The pure rotational band could not be evaluated in this fashion since no measurements for the total band exist. However, mean absorption coefficients were calculated in Ref. 11; these have been included in the present tabulation.

Approximately 80 published emission and absorption spectra were used in the present analysis. The spectral emissivity at intervals of 25 cm<sup>-1</sup>, together with the optical depth, temperature and total pressure served as inputs in a computer program. When the experimentally determined spectra were given with a higher resolution than 25 cm<sup>-1</sup>, the correspondence smoothed by hand. The absorption coefficient was calculated from

$$k(\mathbf{w}) = \frac{1}{8au} \left[ \frac{\mathbf{W}}{\mathbf{d}}(\mathbf{w}) \right]^2 + \left\{ \left( \frac{1}{8au} \left[ \frac{\mathbf{W}}{\mathbf{d}}(\mathbf{w}) \right]^2 \right)^2 + \left[ \frac{1}{u} \frac{\mathbf{W}}{\mathbf{d}}(\mathbf{w}) \right]^2 \right\}^{\frac{1}{2}}$$
(7)

by choosing  $\overline{a}$  such that Eq. (6) was satisfied (see Fig. 1). If

$$\frac{1}{u}\int \ln \frac{1}{1-\varepsilon(w)} dw \approx \alpha , \qquad (8)$$

the spectrum was from a thin gas sample and the absorption coefficients could be determined directly.

Two sources of error in the determination of  $\overline{a}(T)$  and k(w,T) are recognized:

1) Error in the measurements of the state parameters of the gas

or of the spectral emissivity or absorptivity, and

2) Uncertainty in the integrated intensities of the vibrationrotation bands.

An analysis was performed in which reasonable error limits for the state parameters and the integrated intensities were selected. The state parameters T, u, c,  $p_T$ , and the emissivity were varied by  $\pm$  5%, while  $\alpha$  was varied by  $\pm$  20%. It is believed that these variations constitute upper limits, since many experimenters quote smaller error limits. It is found that the uncertainty of  $\pm$  5% in the measured emissivity has the greatest influence on  $\overline{a}(T)$  and k(w,T) ( $\Delta \overline{a} \approx \pm$  100%,  $\Delta k \approx \pm$  20%). An uncertainty of  $\pm$  20% in  $\alpha$  gives  $\Delta \overline{a} \approx \pm$  20% and  $\Delta k \approx \pm$  15%. All other parameters influence  $\overline{a}$  and k by  $\pm$  5% or less.

#### RESULTS

The average absorption coefficients obtained by the method described in the previous section are tabulated in Table II, and are plotted in Fig. 2 for the two extreme temperatures, 300°K and 3000°K. In general, the mean deviation of the absorption coefficients taken from the individual spectra is within ± 20%. As an example, in Fig. 3 k at 1250 cm<sup>-1</sup> is plotted versus temperature. The solid line represents the mean of the individual points. From this curve, the absorption coefficients at 300°, 600°, 1000°, 1500°, 2000°, 2500°, and 3000°K are selected and entered in Table II. In some portions of the spectra, where the absorption coefficient changes rapidly, a much greater spread in the individual absorption coefficients is observed. In Fig. 4 k at 3750 cm<sup>-1</sup>, the Q-branch region on the 2.7-µ band, is plotted versus temperature. The large spread here is probably introduced by small errors in the wavenumber calibration. Also, the absorption coefficients in the spectral regions between the individual bands ("troughs") and in the region w > 7500 cm<sup>-1</sup> are not very certain

because of lack of sufficient data. Since, at the present time, there are no data for temperatures below 2000°K in this region, no absorption coefficients are listed in Table II.

The average line spacings  $\overline{d}(T)$  obtained from the values of  $\overline{a}(T)$ , using Eq. (4), are plotted in Fig. 5. Data from approximately 80 spectra of the 2.7- and 6.3- $\mu$  band are represented. Each point has about a factor of 2 uncertainty. The overall spread of points is mostly within a factor of 2 to 3, but can be as large as a factor of 10 (at 1200°K for instance). The straight line is a least square fit to the points and is given by

$$\bar{d}(T) = \exp[-.00106T + 1.21]$$
 (9)

One point at 2200°K is obtained from the rotational band. The point from Oppenheim<sup>23</sup> at 1200°K is a simple average taken from his measured  $(2\pi\gamma^{0}/d)$  values (see Fig. 6). The majority of points in Fig. 5 are restricted to temperatures < 1800°K because the high-temperature measurements made with the rocket motor in this laboratory were in the thin gas region where Eq. (8) applies.

The two combination bands at 1.87  $\mu$  and 1.38  $\mu$  are much weaker than the fundamental bands. Room temperature absorption spectra of the 1.87- $\mu$  band indicate that the value of  $\overline{d}$  is practically the same as that obtained from the two fundamental bands. Values of  $\overline{d}$  derived from the data presented by Nelson are generally higher than values of  $\overline{d}$  obtained from his data for the two fundamental bands, but the experimental uncertainty in his measurements does not permit any definite conclusions. Our own data were always obtained at low optical depths, so that no  $\overline{d}$  could be obtained. In the absence of further information, it is assumed that the  $\overline{d}(T)$  for the two combination bands is the same as that deduced for the fundamental bands.

The rotational band is very strong and fine structure affects the measurements even at the highest temperatures measured (~ 2700°K).  $^{11}$  However, since the measurements were carried out only up to 22  $\mu$ , no total band strength could be obtained. Thus, the procedure used before for the vibration-rotation bands could not be employed. Additional knowledge about the integrated intensity of this band up to 22  $\mu$  was needed which presupposed the knowledge of the spectral absorption coefficient. The theoretical relationships of Ref. 11 were used here in a semi-empirical fashion by adjusting the effective mean rotational constant slightly to fit the observed spectra in the far wing of the rotational band. The values of k thus determined were used to calculate

1000 cm<sup>-1</sup>

$$\int_{450 \text{ cm}^{-1}} k(\mathbf{w}, \mathbf{T}) d\mathbf{w} = \alpha'(\mathbf{T})$$
(10)

which then replaced Eq. (7).

#### COMPARISON WITH OTHER THEORETICAL AND EXPERIMENTAL RESULTS

A check of the absorption coefficients derived here against independent work is possible for some selected temperatures. Benedict  $^{24,25}$  has obtained absorption coefficients for individual lines in the 6.3- $\mu$  band at  $\approx 300\,^{\circ}$ K and  $\approx 600\,^{\circ}$ K and in the 2.7- $\mu$  band for  $\approx 300\,^{\circ}$ K. We have averaged Benedict's values and compared them with the absorption coefficients taken from this work (see Figs. 7, 8, and 9). Oppenheim's  $^{23}$  detailed measurement of the absorption coefficients in the 2.7- $\mu$  band at 1200°K is compared in Fig. 10. Goldstein's determination of the absorption coefficients of the 1.38-, 1.87-, 2.7-, and 6.3- $\mu$  bands at  $\approx 500\,^{\circ}$ K and 1200°K are compared with the present values in Figs. 11 through 16. One set of experimental data has recently become available. Simmons, Arnold, and

Smith 26 have measured the 2.7- $\mu$  bands between 700° and 1200°K at various optical paths. Since their data were not included in our derivation of the absorption coefficients and band averaged line spacing, they provide us with an independent check for calculating thick gas spectra of the 2.7- $\mu$  bands at around 1000°K. The result of this comparison for some of the spectra is given in Figs. 17-22. Similar agreement is found with the other spectra from Ref. 26.

In a consistency check, some of the measured spectra, which were used to obtain a band-averaged line spacing and average absorption coefficients, are recalculated using the d from Fig. 5 (dashed line) and the k's from Table II. In Fig. 23 we compare Nelson's measurements of the 6.3- $\mu$  band at llll°K with the present representation of k( $\omega$ ) and T. In Figs. 24, 25, and 26, the spectra obtained in this laboratory are compared at 2200°K. In Fig. 27, one spectrum obtained by Burch and Gryvnak<sup>2</sup> of the 2.7- $\mu$  band at 1200°K is compared with the present results. In general the agreement is adequate for the present representation of H<sub>2</sub>O at high temperatures.

#### CONCLUDING REMARKS

The absorption coefficients averaged over 25 cm<sup>-1</sup> intervals given in this study for temperatures greater than 300°K represent the "best" values available at the present time. (The values at 300°K were included here only for completeness. More detailed tabulations are available in the literature, for example, References 24 and 27). The absorption coefficients derived here have been based on nearly all the available published spectra, using a statistical model for the curve of growth and a frequency-independent fine structure term to extrapolate to zero pathlength. In the slowly varying portions of the spectra, the uncer-

tainties are believed to be within ± 20%, while in portions of the spectra where steep changes occur, the uncertainties may be higher. Also the values in the troughs between the vibration-rotation bands are not very certain because of the lack of sufficient data.

The different frequency-independent fine structure terms obtained from the analysis of the optically thick spectra were used to calculate band-averaged line spacings as functions of temperature. These individual points are uncertain by a factor of 2 and although the spread can be as high as a factor of 10 (at 1200°K for instance), the mean values (dashed curve in Fig. 5) are believed to be uncertain by a factor of 2 only. That these mean values may be used for first approximations in the calculation of moderately thick gases was shown in the comparison with some of the spectra by Simmons, et.al., <sup>26</sup> in Figs. 17-22. More accurate calculations for very thick gases will require a detailed evaluation of the dependence of the fine structure parameters on both temperature and wavelength.

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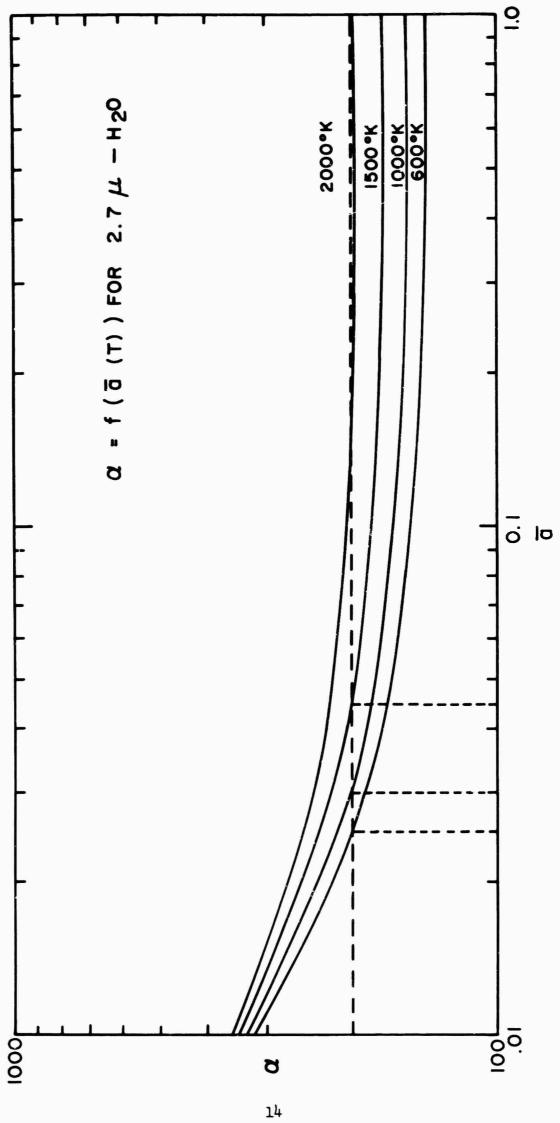


Figure 1. The integrated intensity  $\alpha$  of the 2.7- $\mu$  band as function of  $\bar{a}$  for various temperatures.

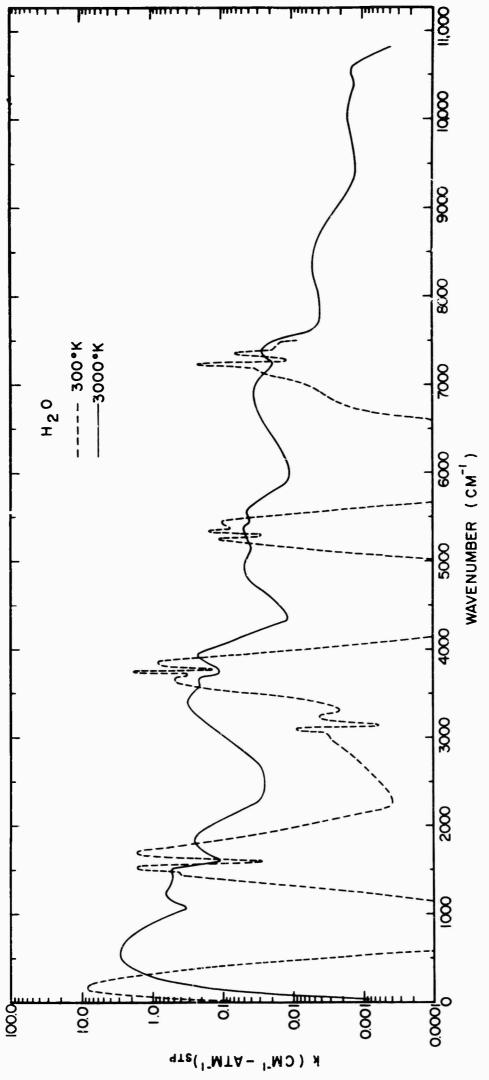


Figure 2. Plot of absorption coefficients at 300°K and 3000°K.

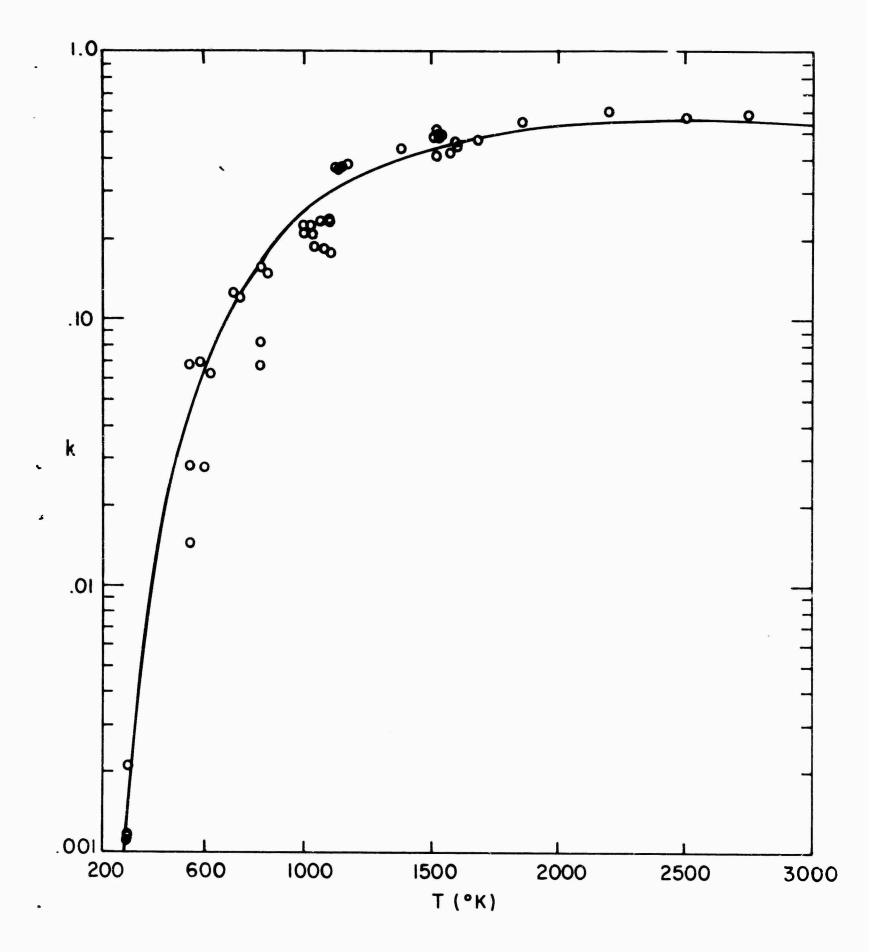


Figure 3. Compilation of average absorption coefficients versus temperature at 1250 cm<sup>-1</sup> as obtained from spectra taken by several investigators. Solid line is the average value

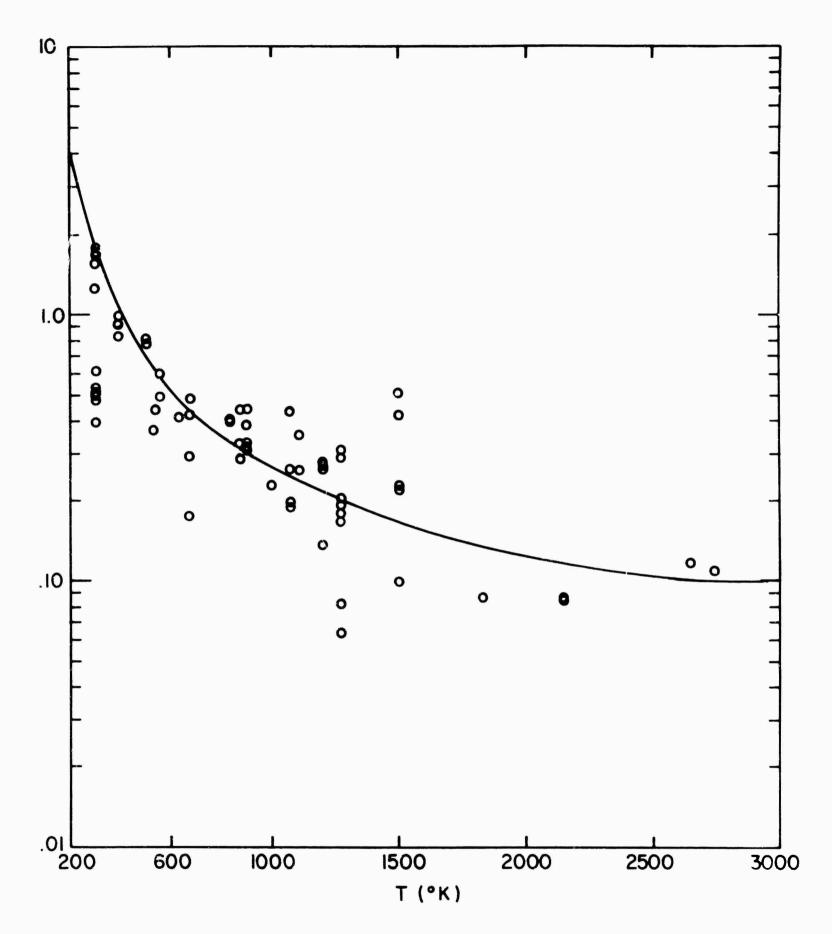
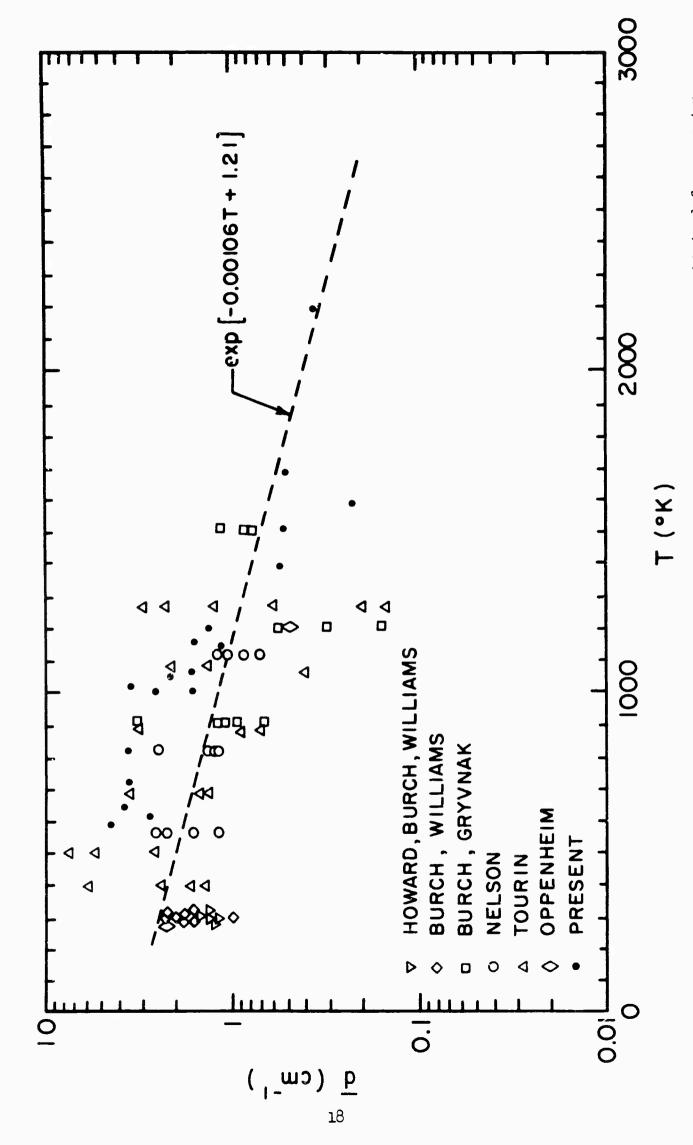
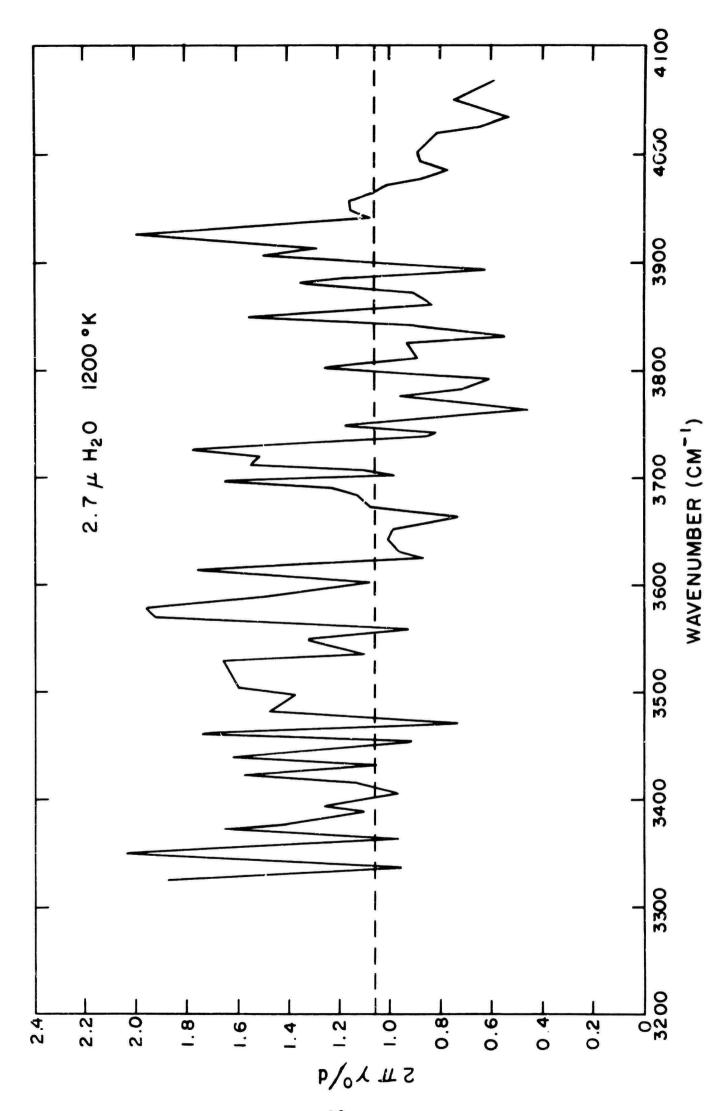


Figure 4. Compilation of average absorption coefficients versus temperature at  $3750~{\rm cm}^{-1}$  as obtained from spectra taken by several investigators. Solid line is the average value.



Compilation of band averaged line spacings d versus temperature as obtained from spectra taken by several investigators. Figure 5.



The ratio of line width to line spacing  $(2\pi\gamma^{0}/d)$  as function of wavenumber, obtained by Oppenheim of the  $2.7-\mu$  band at  $1200^{\circ}K$ . Figure 6.

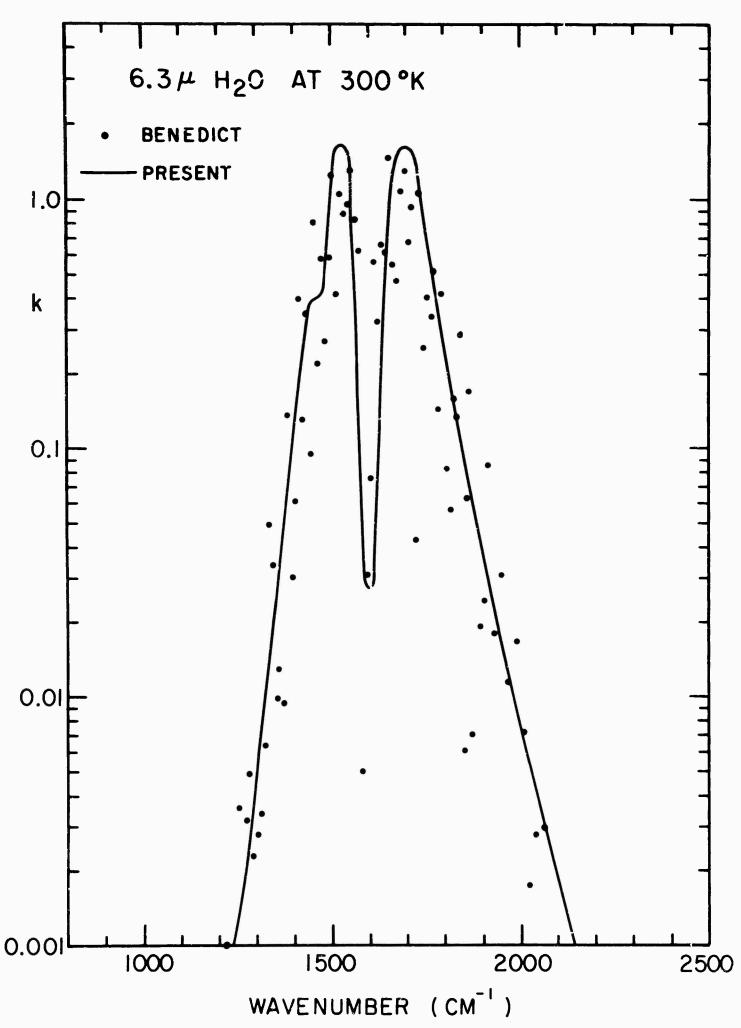


Figure 7. Comparison of absorption coefficients, obtained by Benedict at 300°K, averaged over 25 cm<sup>-1</sup>.

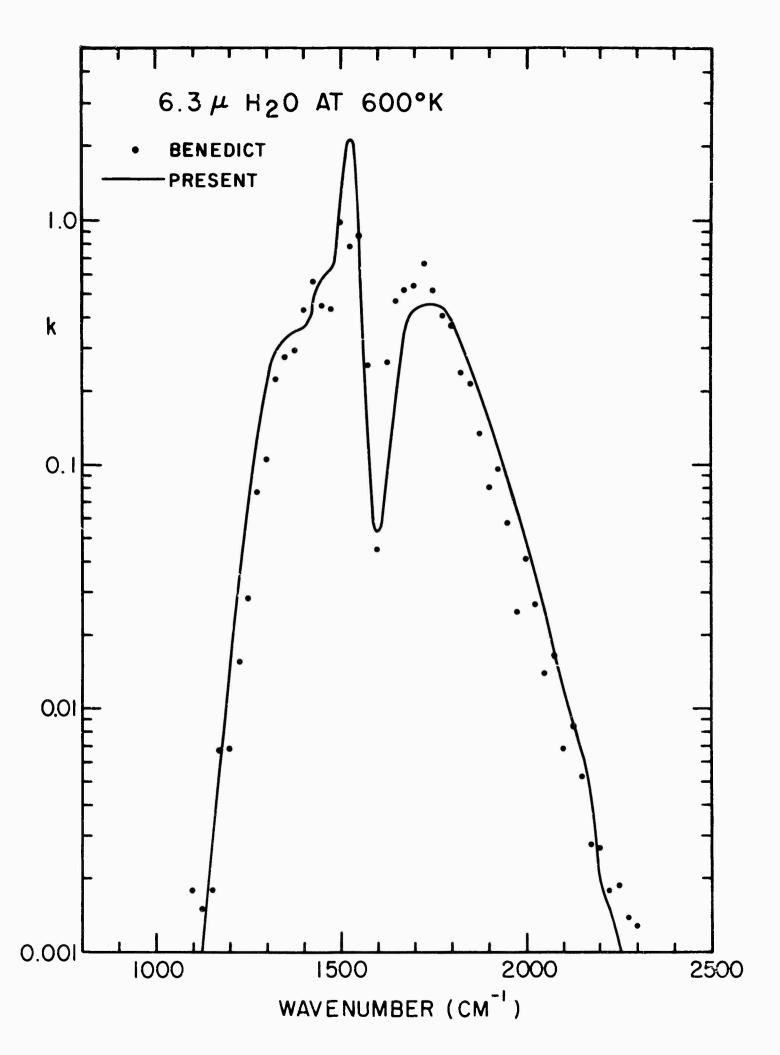


Figure 8. Comparison of absorption coefficients, obtained by Benedict at 600  $^{\bullet}$ K, averaged over 25 cm $^{-1}$ .

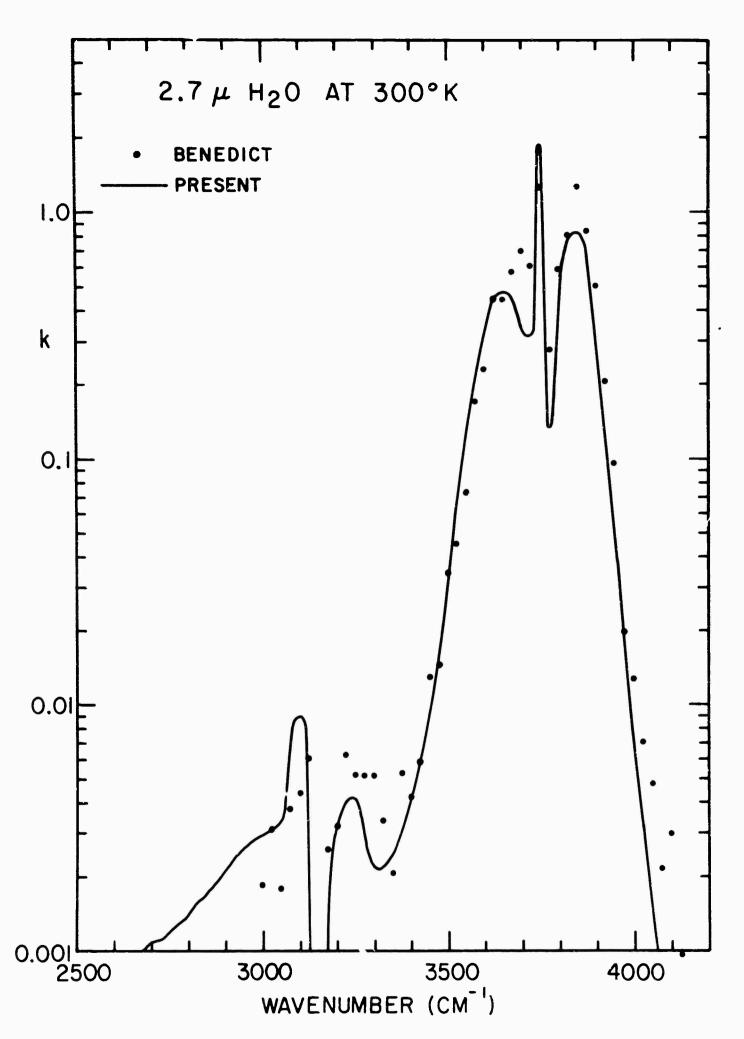


Figure 9. Comparison of absorption coefficients, obtained by Benedict at  $300^{\circ}K$ , averaged over 25 cm<sup>-1</sup>.

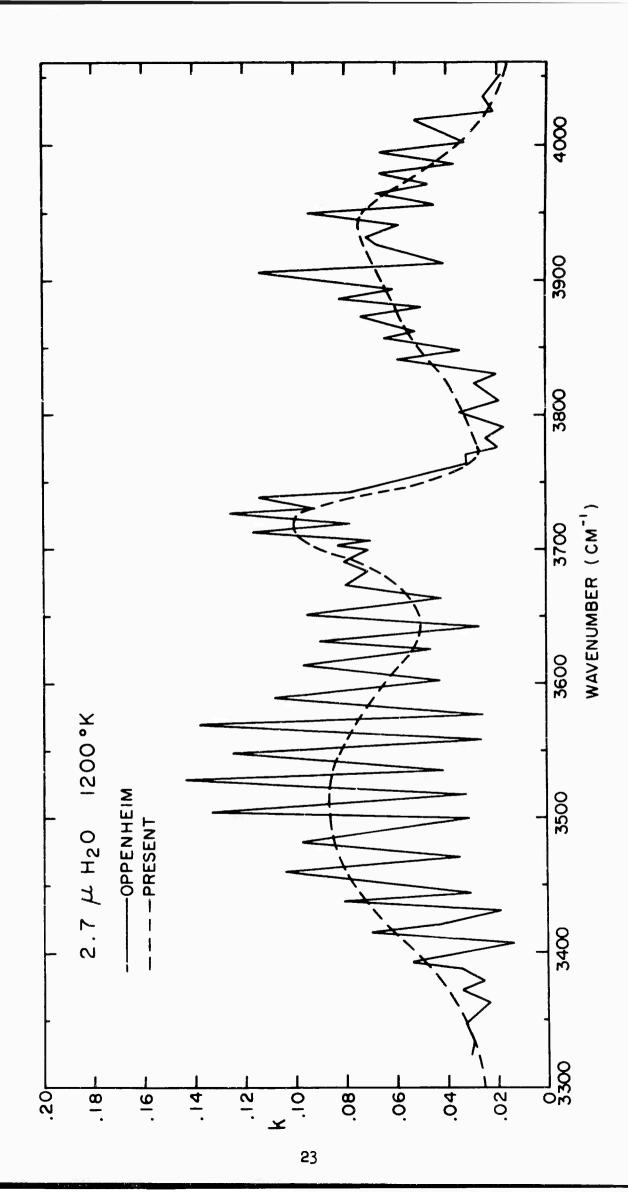


Figure 10. Comparison of absorption coefficients, obtained by Oppenheim.

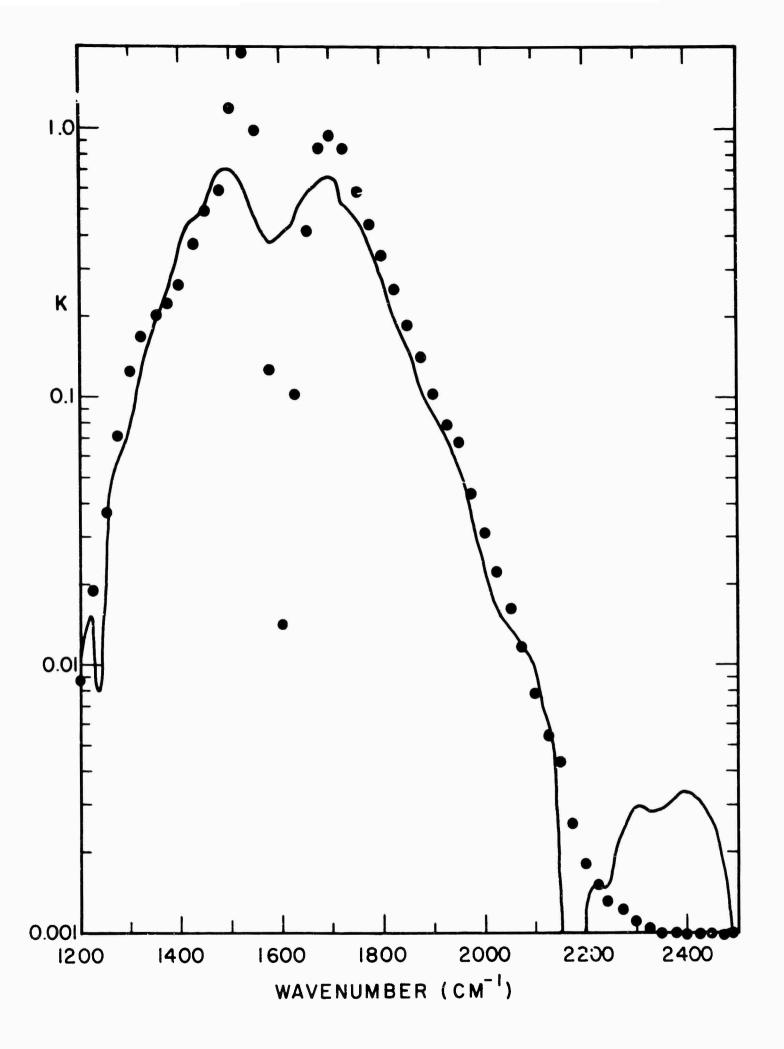


Figure 11. Comparison of absorption coefficients, obtained by Goldstein of the 6.3- $\mu$  band at 473°K (solid line). Present values are given as points.

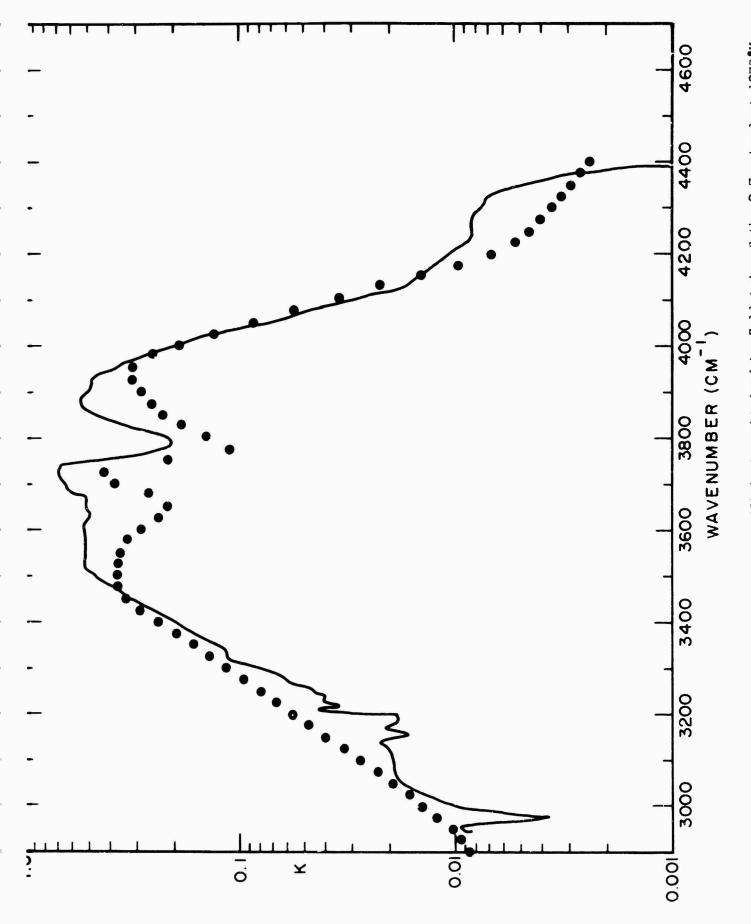


Figure 12. Comparison of absorption coefficients, obtained by Goldstein of the 2.7-w band at 1273°K (solid line). Present values are given as points.

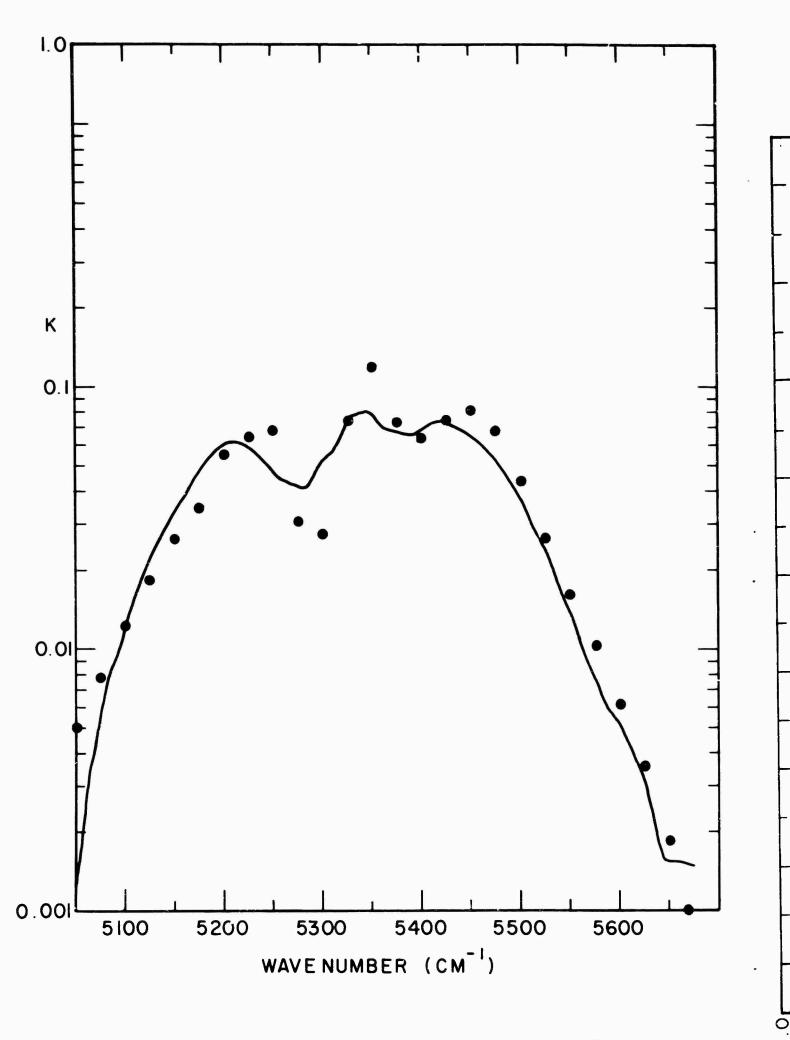
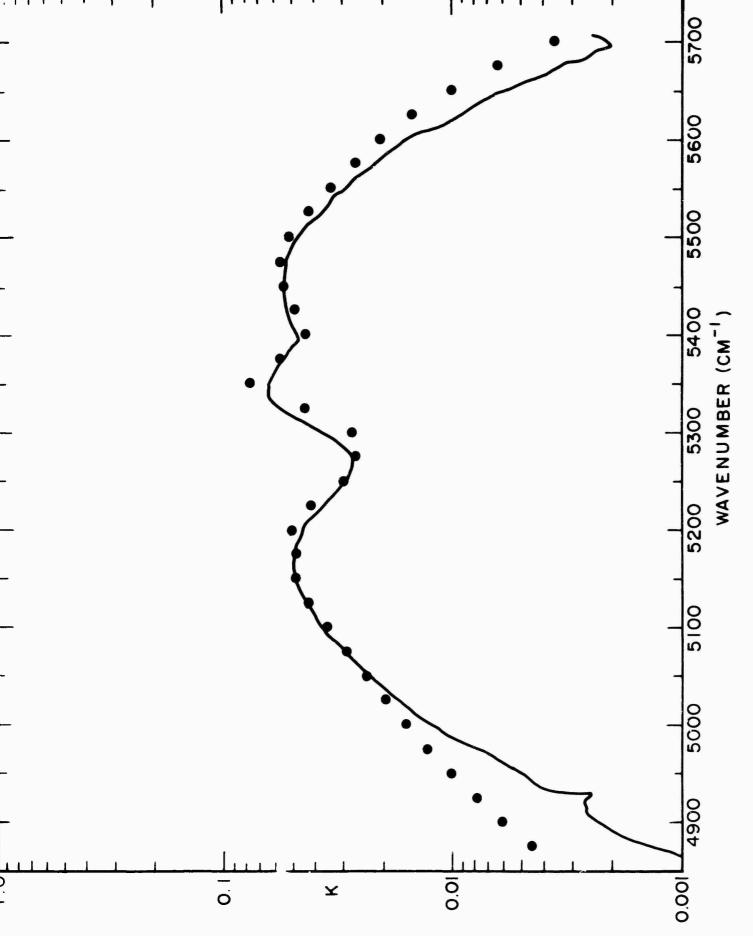
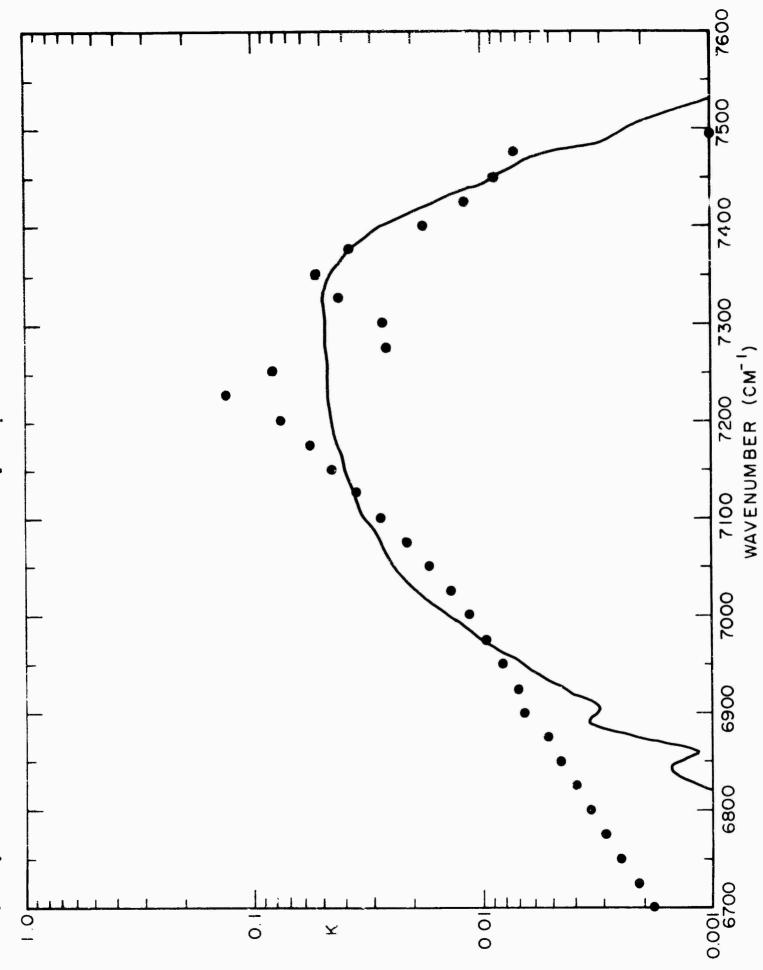


Figure 13. Comparison of absorption coefficients, obtained by Goldstein of the 1.9- $\mu$  band at 473°K (solid line). Present values are given as points.



First No. Cap of emorphion coefficients, obtained by Goldslein of the 1.9-µ band at 373°K (clillis). Present velues are given as points.



. The corplice coefficients, obtained by Goldstein of the  $1.h_{\pm}$ s that is  $1.75^{\circ}{\rm K}$ 

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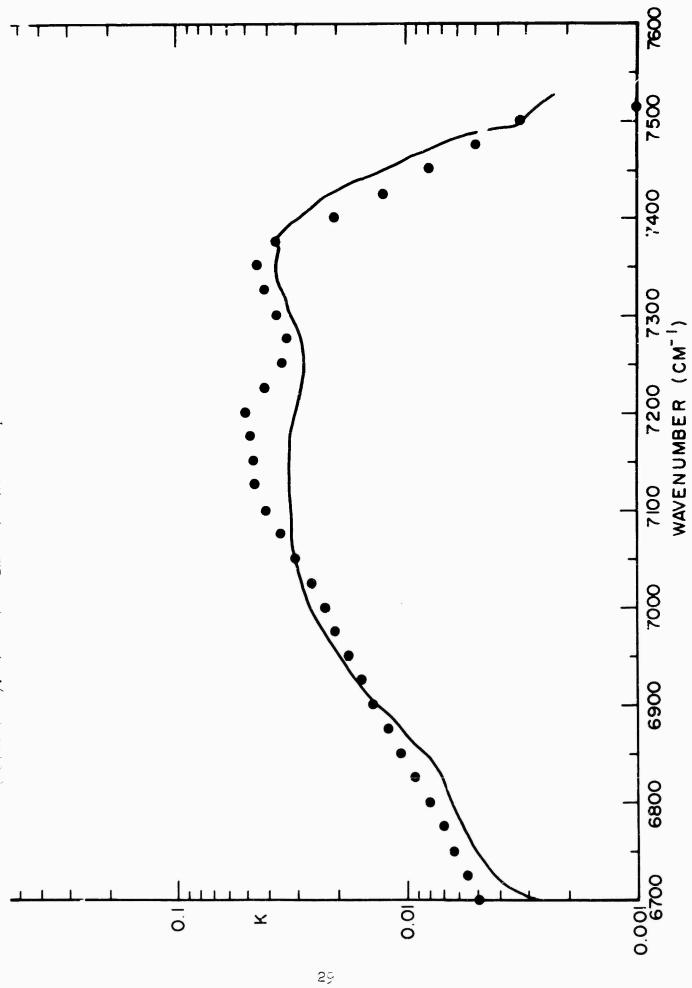


Figure 15. Comperison of absorption coefficients, obtained by Goldstein of the 1.4-m band at 773.7K (solid line). Prerent values are given as points.

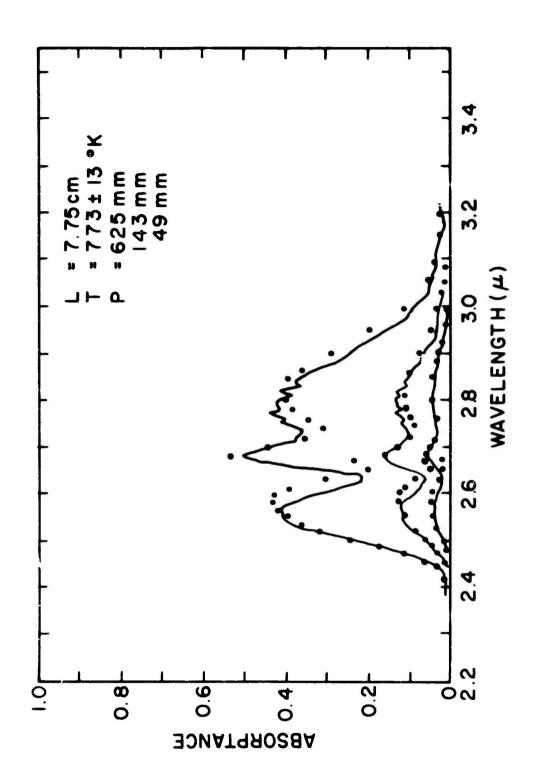
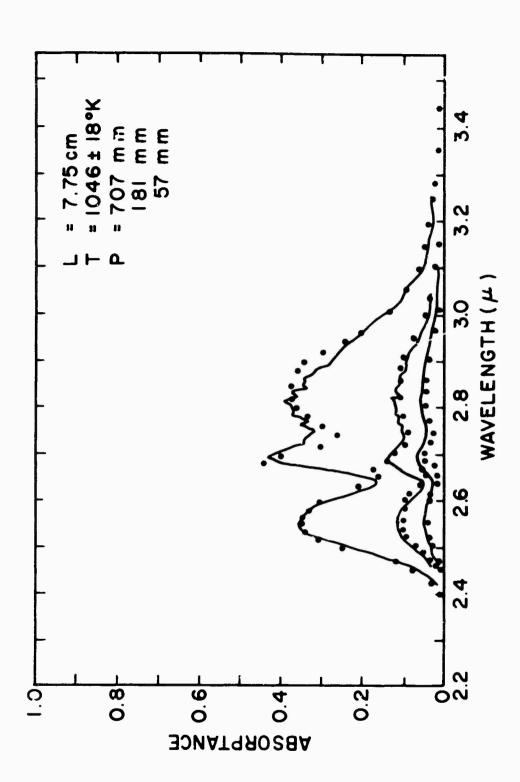
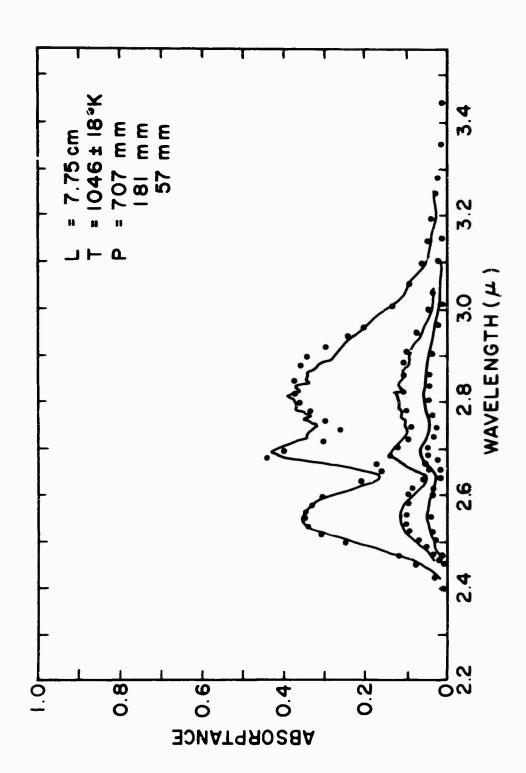


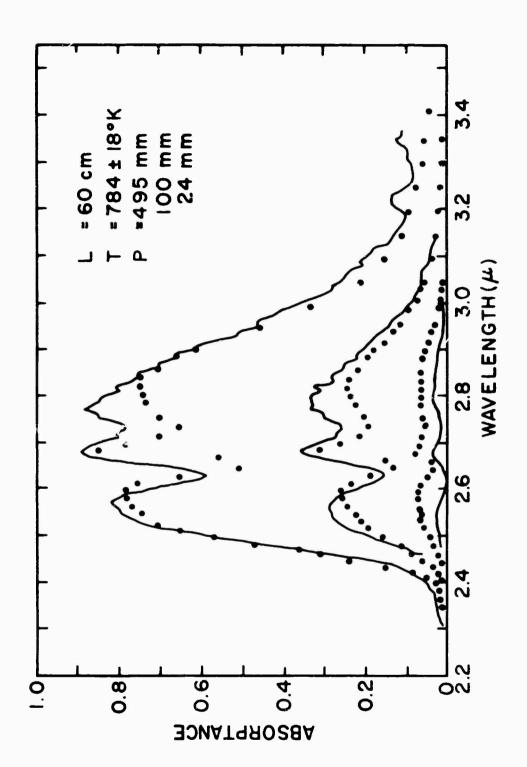
Figure 17. Comparison of measured and calculated emissivity of the 2.7-µ band. Solid lines are the measured spectra taken from Ref. 26, points are the present values.



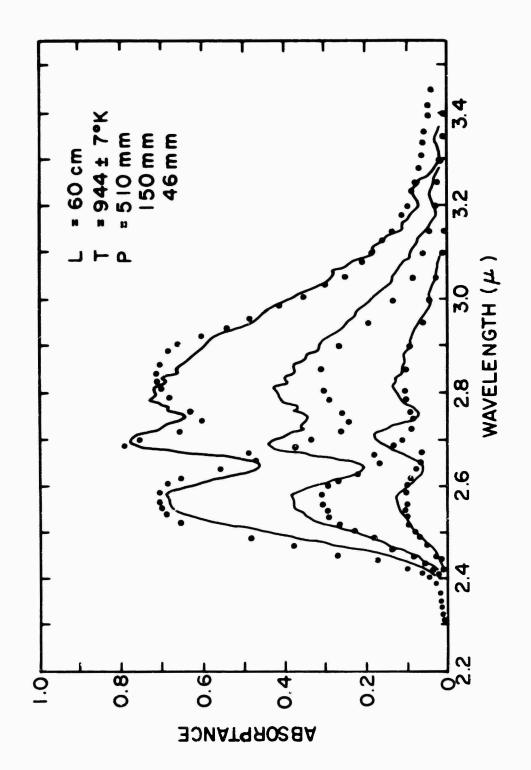
Solid lines are the Figure 18. Comparison of measured and calculated emissivity of the 2.7-4 band. measured spectra taken from Ref. 26, points are the present values.



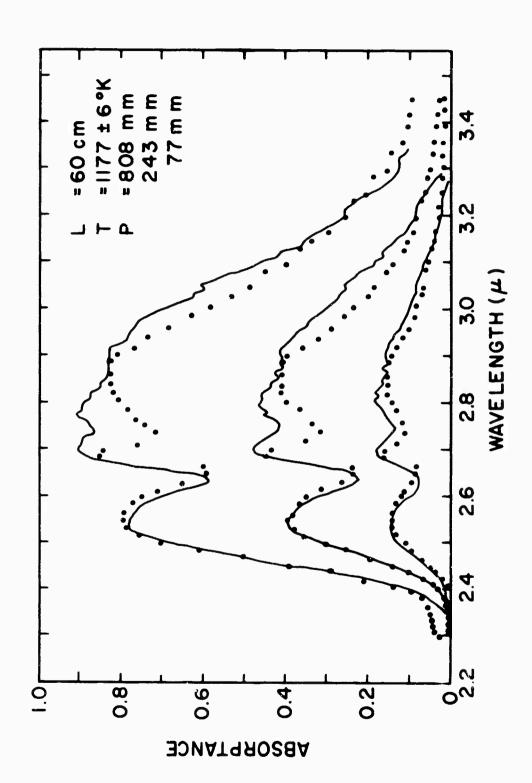
Solid lines are the Comparison of measured and calculated emissivity of the 2.7- band. measured spectra taken from Ref. 26, points are the present values. Figure 19.



Comparison of measured and calculated emissivity of the 2.7-µ band. Solld lines are the measured spectra taken from Ref. 26, points are the present values. Figure 20.



Comparison of measured and calculated emissivity of the 2.7- uband. Solid lines are the measured spectra taken from Ref. 26, points are the present values. Figure 21.



Comparison of measured and calculated emissivity of the 2.7-µ hand. Solid lines are the measured spectra taken from Ref. 26, points are the present values. Figure 22.

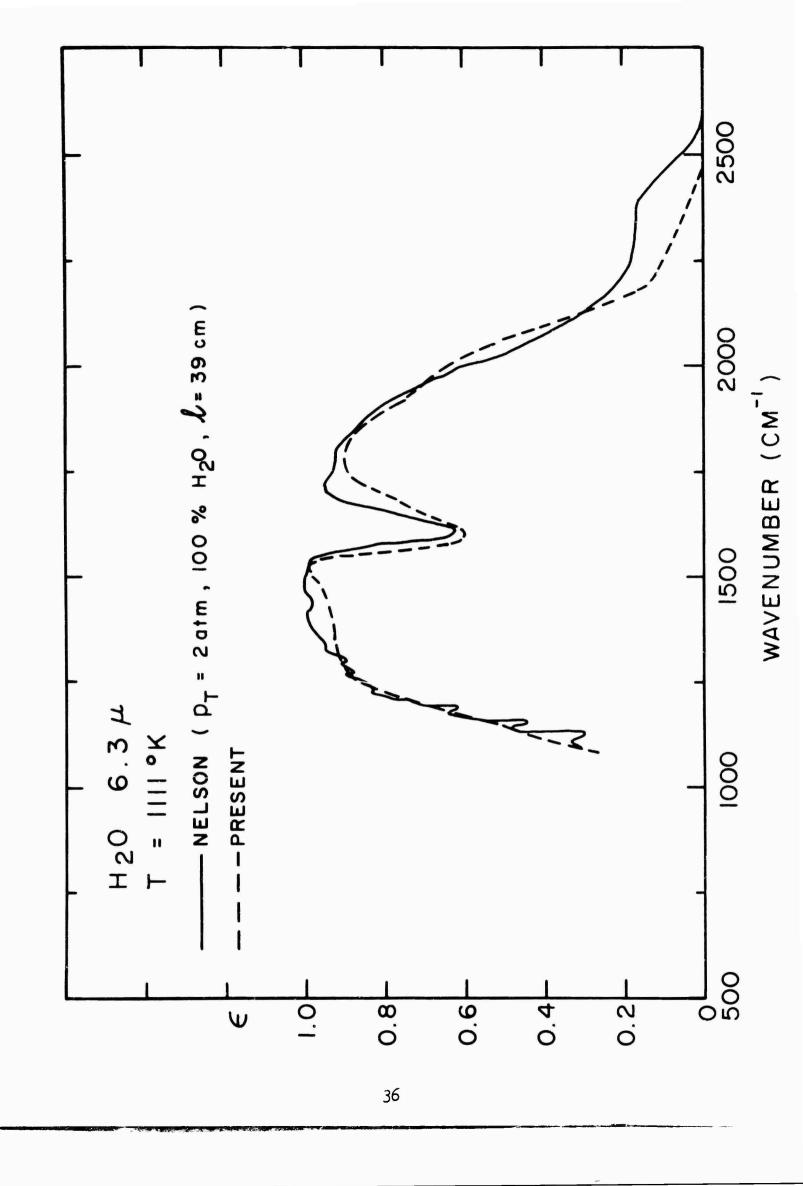
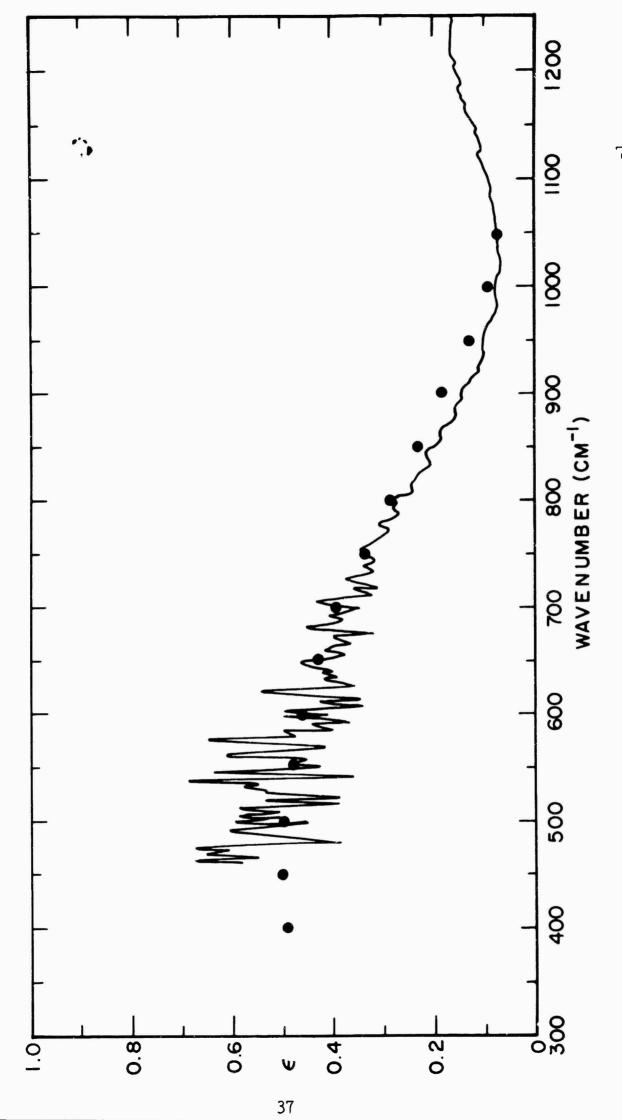
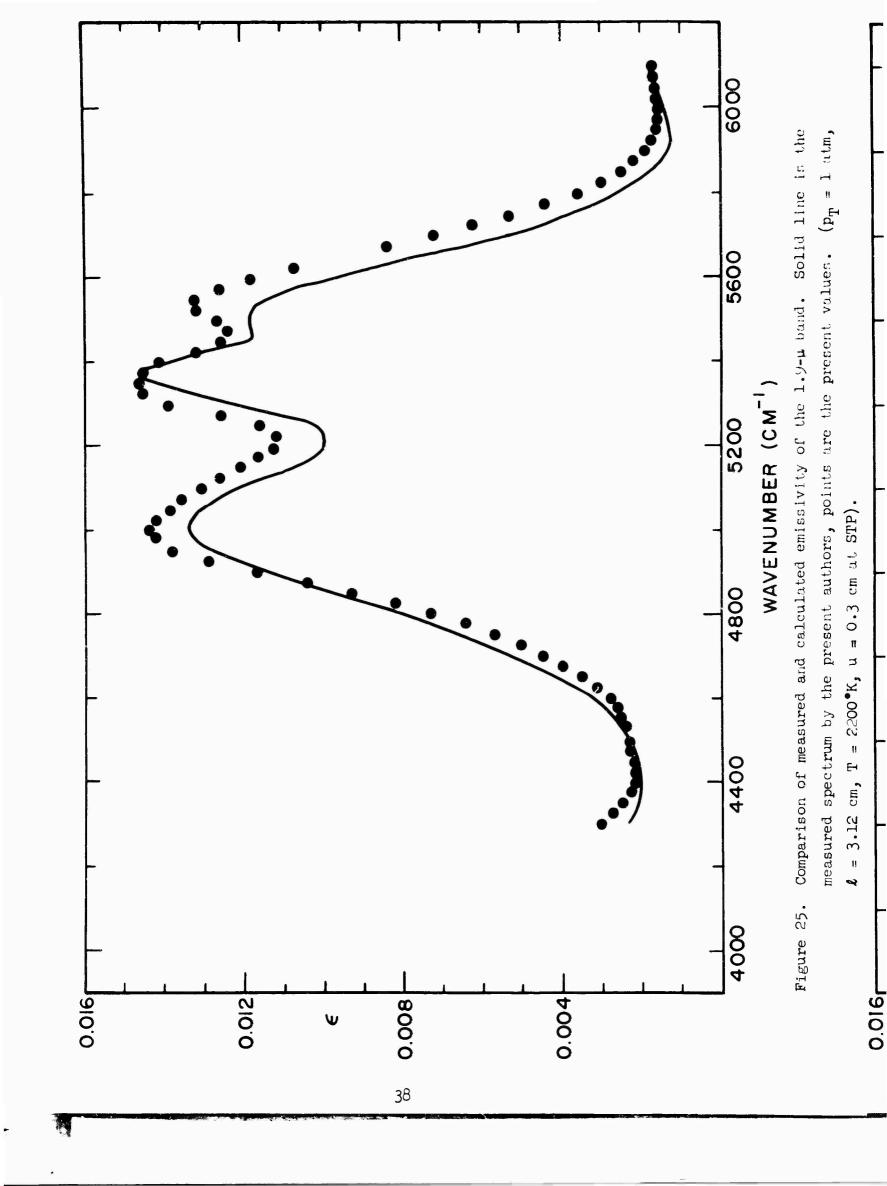
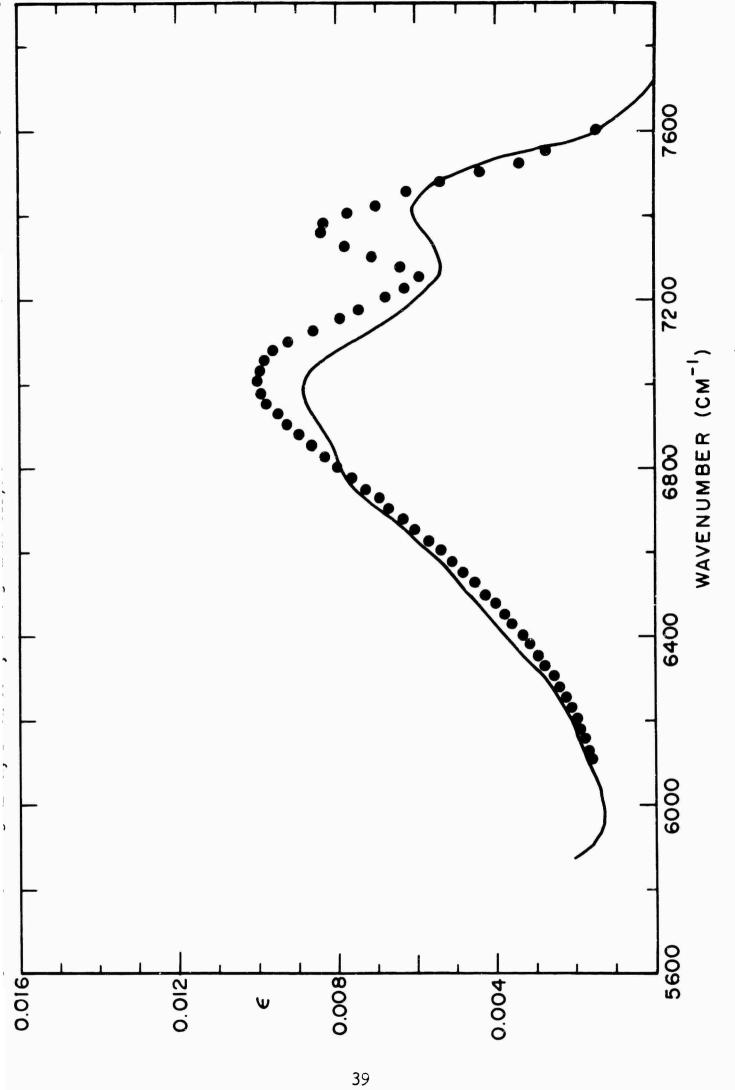


Figure 23. Comparison of measured and calculated emissivity.

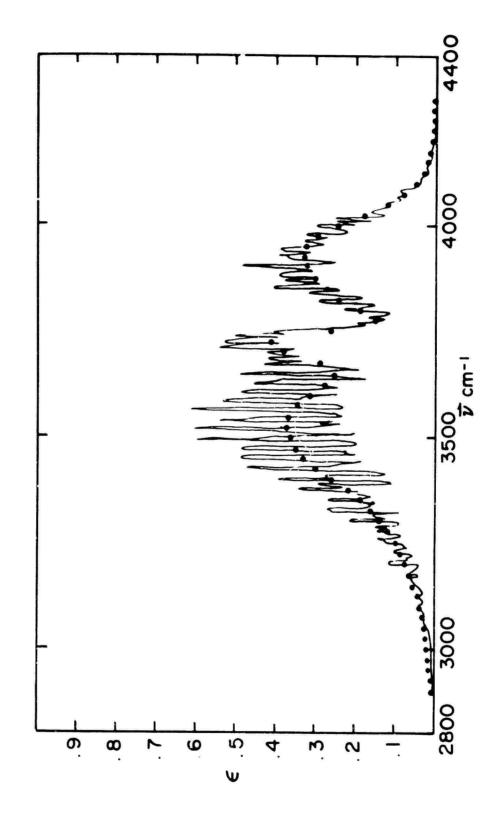


Comparison of measured and calculated emissivity of the rotational band from 500 cm to  $1050~{
m cm}^{-1}$ . Solid line is the measured spectrum by the present authors, points are the present values. ( $p_T$  = 1 atm,  $\ell$  = 3.12 cm, T = 2200°K, u = 0.3 cm at STP). Figure 24.





the measured spectrum by the present authors, points are the present values.  $(p_{\rm T}$  = 1 atm, Comparison of measured and calculated emissivity of the 1.4- µ band. Solid line is 1 = 3.12 cm, T = 2200 K, u = 0.3 cm at STP). Figure 26.



Comparison of measured and calculated emissivity of the 2.7-µ hand. Solid line is the measured spectrum by Burch and Gryvnak. ( $p_T$  = 1 atm, T = 1200°K, u = 1.76 cm at STP, 100% H<sub>2</sub>0). Figure 27.

TI TIBLE

ABSORPTION COEFFICIENTS OF HZO

K (PER CM AT STP)

3000 K	190E-0	670E-0	211E-0	529E-0	109E 0	193E 0	09E 0	454E 0	620E 0	822E 0	104E 0	128E 0	154E 0	182E 0	207E 0	229E 0	24BE 0	.262E 01	273E 0	BOE O	83E 0	282E 0	77E 0	269E 0	258E 0	45E 0	29E 0	13E 0
2500 K	300E-0	134E-0	365E-0	845E-0	168E 0	289E 0	460E 0	679E 0	938E 0	122E 0	149E 0	179E 0	208E 0	237% 0	260E 0	280E 0	295E 0	.304E 01	309E 0	307E 0	301E 0	289E 0	275E 0	260E 0	241E 0	221E 0	SOCE O	178E 0
2000 K	450E-0	195E-0	77E-0	128E 0	260E 0	450E 0	695E 0	101E 0	135E 0	72E 0	213E 0	249E 0	284E 0	312E 0	330E 0	341E 0	345E 0	.342E 01	334E 0	319E 0	300E 0	276E 0	248E 0	22E 0	195E 0	69E 0	46E 0	24E 0
1500 K	114E-0	375E-0	104E 0	214E 0	415E 0	765E 0	26E 0	179E 0	230E 0	281E 0	328E 0	361E 0	383E 0	394E 0	396E 0	388E 0	370E 0	.343E 01	310E 0	274E 0	238E 0	204E 0	174E 0	145E 0	118E 0	950E 0	48E 0	BOE O
1 000 K	420E-0	113E 0	300E	650E 0	121E 0	189E 0	261E 0	337E 0	407E 0	456E 0	479E 0	484E 0	471E 0	443E 0	400E 0	347E 0	292E 0	.236E 01	188E 0	145E 0	110E 0	818E 0	<b>598E</b> 0	427E 0	294E 0	200E 0	134E 0	UZE-O
<b>6</b> 00 <b>K</b>	03E 0	65E 0	90E 0	O1E O	25E 0	52E 0	40E 0	OOE O	22E 0	92E 0	28E 0	SOE 0	70E 0	89E 0	OSEO	43E 0	SOE 0	• 61 CE 00	86E 0	36E 0	44E 0	20E-0	45E-0	42E-0	27E-0	30E-0	00E-0	0=30
300 K	50E 0	08E 0	86E 0	SOE 0	25E 0	70E 0	10E 0	82E 0	93E 0	16E 0	99E 0	13E 0	85E 0	93E 0	38E C	20E-0	55E-0	.940E-02	40E-0	05E-0	50E-0	<b>26E-0</b>	30E-0	S0E-0	10E-0	70E-0	70E-0	95E-0
1 / CM	•	ູດ	000	25.	50.	75.	• 00	25.	50.	75.	• 00	25.	50.	75.	00	25.	50.	475.0	00	25.	50.	75.	000	25.	50.	75.	000	25.

554E **537E 518E 578E** SZOE 514E . 159E 516E .177E .143E .955E • 600E .358E •449E **330E** . 126E . 111E • 825E 705E .510E .425E .329E .357E .507E .568E . 604E . 632E .637E • 60 BE 000 0000 00 00 000 00 00 00 00 00 00 ¥ 2500 **365E** .283E .380E SOUE **560E** .312E .314E .450E .580E .570E .535E .515E 490E **530E** \* 156E .136E • 117E .595E .492E •405E .352E **• 289€** .281E .505E .555E • 100E •855E .480E .483E .7186 0000 00 00 00 000 00 00 00 00 00 00 00 00 (CONT.D) ¥ •190E 2000 **560E** 278E 335E **510E** 520E \$10E 8.45E 239E 522E SZOE **560E 595E** •103E .695E .450E 357E 211E • 195E .191E • 194E .202E .223E .260E .410E .470E •497E 498E .485E •495E 000 00 00 0 00 00 00 0 000 00 00 00 .865E-01 730E-01 745E-01 .830E-01 .665E-01 .630E-01 .620E-01 .630E-01 .675E-01 ¥ .122E COEFFICIENTS OF CM AT STP) •228E .630E .174E • 100E .330E .242E .123E .170E .428E • 488E .497E •495E .489E .498E .291E .355E .463E .540E .584E 8 00 0 00 0 00 00 .166E-01 .167E-01 .175E-01 .188E-01 .208E-01 .233E-01 .268E-01 .345E-01 .555E-01 450E-01 .355E-01 .289E-01 .245E-01 .214E-01 .189E-01 .174E-01 .165E-01 .900E-01 290E-0 •152E •249E .475E •435E .340E 390E .458E 462E 505E **500E** .660E ABSORPTION K (PER .950E-03 00 00 00 00 00 620E-03 2755-03 .113E-03 500E-04 230E-04 .103E-04 460E-05 2056-05 .140E-05 350E-05 850E-05 215E-04 570E-04 .150E-03 .380E-03 .245E-02 620E-02 .148E-01 330E-01 635E-01 ¥ 900 212E 361E 285E 328E 345E 460E .123E 569E **560E-06** .275E-05 470E-05 .810E-05 •136E-04 .235E-04 .400E-04 .680E-04 .120E-03 200E-03 .680E-03 500E-02 .385E-07 .670E-07 .113E-06 .195E-06 328E-06 .950E-06 .160E-C5 365E-03 .130E-02 .250E-02 219E-01 .485E-01 •680E-07 .103E-01 ¥ 300 .114E 249E • 418E .397E 100.0 275.0 325.0 350.0 375.0 825.0 950.0 050.0 1125.0 1150.0 200.0 225.0 250.0 300.0 400.0 750.0 775.0 800.0 850.0 875.0 0.006 925.0 975.0 0.000 0250 0.270 1175.0

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		₹	BSORPTION K (PER	COEFFICE	ENTS OF	180	(CONT.D)				
1./CM	3000	¥	У О	x 000	1500	¥	2000 K	2500	¥	3000	¥
00	08E	21	25E 0	15E 0	.894	00	719E 0	OBE	00	25E	8
25.	<b>65E</b>	7:	11E 0	35E 0	.757	00	349E 0	<b>66E</b>	00	30E	00
50.	42E	10	75E 0	4COE O	.310	00	BOE O	<b>70E</b>	00	<b>77E</b>	00
75.	3E	00	35E 0	21E O	.122	00	140E 0	60E	00	73E	0
600	75E-	11	38E-0	69E-0	.938	.0.	04E 0	10E	00	10E	00
625.	23E	00	105-0	935E-0	.101	00	60E-0	360	00	1 1 E	00
650.	30E	00	0 366	35E 0	. 120	00	0 36 O	ZOE	00	35E	00
675.	49E	21	46E 0	79E 0	.151	00	30E 0	40E	00	SÓE	0
700.	62E	21	33E 0	26E 0	.192	00	60E 0	64E	00	73E	00
725.	35E	10	58E 0	90E 0	.237	00	OIE O	30E	00	91E	00
750.	40E	00	67E 0	70E 0	• < 84	00	40E 0	18E	00	360	00
775.	ZOE	00	43E 0	COEO	.306	00	75E 0	40E	00	26E	00
800.	SOE	00	93E 0	O BOO	,312	00	O BOU	<b>56E</b>	00	42E	00
825.	47E	00	19E 0	SOE 0	• 300	00	291E 0	62E	00	SOE	00
850.	10E-	10	52E 0	93E 0	.284	00	76E 0	<b>56E</b>	00	47E	00
875.	BOE-	10	98E 0	44E 0	• 252	00	248E 0	45E	00	38E	00
.006	70E-	21	53E 0	OOE O	• 250	00	SSSE 0	225E	00	22E	00
925.	44E-	10	18E 0	62E 0	• 192	00	OOE O	05E	00	06E	00
950.	62E-	21	00E-0	33E 0	.166	00	180E 0	188E	00	BSE	00
975.	12E-	11	65E-0	07E 0	.142	00	158E 0	<b>62E</b>	00	<b>66</b> E	8
.000	80E-	25	90E-0	870E-0	.121	00	138E 0	43日	00	47E	00
2025.0	.540E-0	02	•352E-01	67	•981E	01	•116E 00	•123E	00	• 129E	00
020	<b>BOE-</b>	25	52E-0	515E-0	• 790	0	0-3066	OZE	00	1 1E	00
075.	<b>60E</b> -	25	79E-0	90E-0	. 60SE	0	830E-0	70E-	01	60E-	70
100.	<b>BOE</b> -	25	23E-0	290E-0	•471E	0	700E-0	30E-	01	40E-	70
125.	27E-	25	30E-0	13E-0	.359E	0	85E-0	30E-	01	25E~	70
150.	80E-	23	80E-0	155E-0	.268E	0	80E-0	30E-	01	628E-	2
175.	20E-	93	00E-0	13E-0	.201	0	395E-0	<b>60E</b> -	01	30E-	10
200-	BOE-	<b>6</b> 0	90E-0	10E-0	•161E		325E-0	95E-	01	556-	01
25.	05E-	23	40E-0	20E-0	•129E	0	67E-0	40E-	0	98E~	010

	3000	S	1 ZE	• 290E-(	75E	67E	365	S		36 t	47	48E	49E	3	53E	58E	63E	<b>70E</b>	BOE	95E	1 8E	43E	78	176	50	0	50	O SE	•665E-(	Ö	.801E-
	2500 K	92E-	53E-	7E	-360	91E-	74-	<b>BE</b> -	•164E-01	63E-	65E-	7E-	72E-	79E-	87E-	96E-	05E-	14E-	29E-	48E-	68E-	82E-	.308E-01	35E-	65E	05E-	50E-	OF	រហ	OEL	00
(CONT.D)	2000 K	18E-	81E-	OE-	45E-	•132E-01	5E-	18E-	11E-	SE-	02E-	00E-	-300	0	02E-	05E~	08E-	7.4	20E-	26	35E-	49E-	•169E-01	91E-	12E-	40E-	72E-	•305E-01	48E-	98	.450E-01
NTS OF H20 P)	1500 K	03E-0	10E-0	0E-0	60E-0	20E-0	90E-0	70E-0	0 E-	50E-0	50E-0	40E-0	40E-0	40E-0	40E-0	30E-0	60E-0	70E-0	80E-0	10E-0	40E-0	60E-0	00E-0	60E-0	SOE-0	00E-0	-36	2E-		76-	.181E-01
ŭ	0	20E-0	40E-0	80E-0	40E-0	00E-0	BOE-0	60E-0	0E-	40E-0	30E-0	30E-0	30E-0	30E-0	30E-0	40E-0	50E-0	50E-0	60E-0	70E-0	80E-0	90E-0	00E-0	1 0E-0	30E-0	50E-0	70E-0	00E-0	40E-0	80E-0	40E-0
ABSORPTION K (PE	<b>6</b> 00 K	00E-0	E-0	60E-0	E-0	E-0	0=30	0E-0	CE-	30E-0	0E-0	P	0-30	E-0	30E-0	30E-0	30E-0	30E-0	40E-0	40E-0	50E-0	50E-0	P	30E-0	0-306	0-300	10E-0	30E-0	50E-0	10E-0	30E-0
	300 K	86E-0	85E-0	90E-0	00E-0	15E-0	30E-0	55E-0	.485E-03	15E-0	50E-0	95E-0	40E-0	80E-0	40E-0	95E-0	55E-0	10E-0	90E-0	06E-0	13E-0	22E-0	30E-0	45E-0	9E-0	74E-0	95E-0	13E-0	8E-0	60E-0	82E-0
	1/CM	250.	275.	300	325	350.	375.	4004	2425.0	450	475.	500	525.	550.	575.	600	625.	650.	675.	700	725.	750.	775.	800	825.	850.	875.	006	925.	950	975

385E-01 .970E-01 • 106E •117E •129E . 142E .170E • 187E .293E .302E • 30BE • 300E . 290E .272E .265E .241E • 222E .208E .201E • 155E • 202E • 220E .241E •262E .280E .210E .213E .200E • 150E . 1 1 3E 00 00 00 00 00 00 8 00 00 00 00 00 9 00 00 00 00 .104E 00 .855E-01 .945E-01 .278E .291E .286E .117E . 129E •160E .219E .240E •253E .285E .291E .293E .270E .255E .237E .215E .202E •144E .178E •199E .265E . 196E 198E .197E .172E 00 00 00 00 00 00 00 00 00 995E-01 HEO CONT.C 510E-01 590E-01 .760E-01 880E-01 •665E-01 •112E 305E .128E • 185E .225E •147E •166E .205E .226E .270E -302E .312E .312E -264E .245E .200E • 188E .229E •175E -246E .288E .284E . 188E .207E 00 00 00 00 00 00 00 00 00 0 00 00 00 00 00 00 00 00 .420E-01 .498E-01 .608E-01 .725E-01 .870E-01 211E-01 .247E-01 290E-01 .347E-01 Y •139E .101E •119E ABSORPTION COEFFICIENTS OF .162E .214E .341E • 328E • 187E •245E .278E .311E .350E .343E .309E .279E .243E .216E .206E .250E .3;9E .295E K (PER CM AT STP) ၀ 00 00 00 00 00 00 00 00 8 00 .660E-02 .760E-02 00 •590E-02 .870E-02 • 1 OCE-01 .117E-01 .138E-01 .167E-01 .202E-01 .252E-01 .321E-01 .410E-01 .540E-01 .710E-01 .980E-01 •135E •187E •330E SBOE .251E .370E .400E .400E .390E •378E •260**€** .225E •330E -270E 460E 00 00 00 0 00 00 00 .370E-02 .400E-02 .450E-02 480E-02 510E-02 600E-02 .700E-02 00 .330E-02 550E-02 .860E-02 .103E-01 .129E-01 .161E-01 .212E-01 .285E-01 .385E-01 540E-01 .770E-01 600 • 117E •173E 375E .258E .401E .401E 392E .320E 255E 340F 518E 00 00 340E-02 .100E-02 640E-03 330E-02 220E-02 220E-02 25CE-02 310E-02 420E-02 360E-01 295E-02 31CE-02 730E-02 900E-02 .160E-02 410E-02 410E-02 290E-02 600E-02 .940E-02 .720E-01 •133E.00 .165E-01 215E 318E .442E 448E .473E 328E 3100.0 3125.0 3150.0 3175.0 3200.0 3225.0 3275.0 3300.0 3325.0 3350.0 3375.0 3400.0 3425.0 3450.0 3475.0 3500.0 3525.0 3550.0 3575.0 3600.0 3625.0 3025.0 3050.0 3075.0 3250.0 3650.0 3675.0 3700.0

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H20	
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ON CO	
ABSORPTION COEFFICIENTS OF HZO (CONT.	

		ABSORPTION K (PE	COEFFICIER CM AT ST	NTS OF HZO	(CONT.D)		
1/CM	300 K	¥	1000 K	1500 K	2000 K	2500 K	3000 K
500	00E-0	10E-0	COE-0	20E .0	0-300		1 E-
525.	OOE-C	OEL	.9005-03	30E-0	20E-0	36-	Ĭ W
550.	0-300	30E-0	0E-0	30E-0	SOE-0	8E*	73E-
575.	00E-0	105-0	COE-0	40E-0	90E-0	•115E-01	7E-
600	00E-0	30	•900E-03	E-0	40E-	らたー	0
625.	OCE-0	70E-	OE-C	80E-0	10E-0	39	Ū
650.	COE-C	0-300	0E-0	10E-0	-300	.162E-01	Ī
4675.0	•300E-09	.210E-06	.11CE-02	50E-	• 104E-01	9	70E-
700	00E-0	30E-	0=30	30E-0	19E-	J)	.300E-01
725.	00E-0	80E-0	0=30	30E-0	32E-	9	38E-
750.	00E-0	755-	0-30	0E-0	4	.292E-01	1
775.	00E-0	70E-0	10E-0	055-0	70E-	22	. 399£
800	50E-0	50E-0	0E-0	36E-	00E-	52	22E-
825.	30E-0	53E-0	70E-0	70E-	33E-	96	04
850.	65E-0	155-0	00E-0	0-370	73E-	03E-	58E-
875.	15E-0	30E-0	0E-0	416-	20E-	S	65E-
900	0 /E-0	45E-0	80E-0	10	75E-	50	Ī
925.	75E-0	90E-0	13E-0	12E-	20E-	20	78E-
950.	70E-0	-300	45E-0	48E-	59E-	85	78E-
975.	65E-0	45E-0	80	81E-	BOE-	8	3E-
000	20E-0	90	OE-	17E-	9		• 460E-0
025.	15E-0	30E-0		47	9	72	Ü
650.	20E-0	60E-	08E-	80	80E-	55E-	0E=
075.	20E-0	30E-	Ø	00E-	70E-	35E-	15E-
100.	00E-0	98	ı	1	55E-	8	Ī
125.	30E-0			• 490E-01	ı	02	90E-
150.	00E-0	-306	.524E-01		188	92	n
175.	80E-0	62E-	BOE-	75E-	95E-	9E-	93E-
00	48E-0	Ŵ	OE	-30E	75E-	2E-	00
225.	18E-0	06	.330E-01		•369E-01	•400E-01	.418E-01

	3000	4	30E-	H	~	F	.478E-0	60E-	25E-	8	05E-	1 8E-	38E-	4	30E-	05E-	Ļ	30E-	•320E-0	HAE-	58E-	Ī	07E-	• 184E-0	36-	45E-	30E-	OF	16	12E-	• 109E-0
	2500 K	.1.	96-		BOE-		80€−	70E-	38E-0	18	2E-0	i W	38E-	35E-0	20E-0	95E-0	62E-0	30E-0	0E-	69E-0	39E-0	=======================================	83E-0	7E-0	33E-	10E-0	0E-0	30E-0	0E-0	ı	0E-0
(CONT.C)	2000 K	0	20	5E-	.498E-01	Õ		.485E-01	50E-	OE-	SOE-		55E-0	556-	28E-0	00E-	60E-	12E-	3E-0	30E-	95E-0	63E-0	33E-0	E-0	20E-	008-0	1 0E-0	40E-0	80E-	0E-0	0
NTS OF H20 P)	1500 K	30	-309	L	65	8E-	53	.523E-01	93E-	45E-	Ĭ	50E-	68	48E-	10E-	63E-	17E-	59E-	11E-0	68E-	32E-	E-0	40E-0	50E-0	20E-	E-0	80E-0	30E-0	00E-0	0E-0	E-0
COEFFICIE R CM AT ST	100		50E-	70	•440E-01	.630E-01	-306	C	39E-	72E-	02E-	.502E-01	55E-	BOE-	10E-	50E-	98E-0	31E-0	60E-	00E-0	90E-0	OE-D	00E-	0-30C	-300	-300	.500E-03		.500E-03	i W	
ABSORPTION K (PE		68E-	B)	270E-	22	• 105E 00	710E-0	483	SET	682E-	680E-0	ហ	350E-0	38E-0	158E-	101E-0	590E-0	310E-0	130E-0	00E-0	55	•145E-04	68E-0	5	9	ဂ္ဂ	88E-	30	ល	0-30	• 495E-04
	300 K	11E 0	29E-0	B1E-0	21E 0	39E 0	74E-0	58E-0	.985E-01	96E-0	90E-0	25E-0	50E-0	20E-0	70E-5	13E-0	60E-0	85E-0	660E-0	70E-0	10E-0	55E-0	98E-0	70E-0	25E-0	30E-0	80E-0	30E-0	00E-0	005-0	50E-0
	1 / CM	250.	275.	300	325.	350.	375.	400	5425.0	450.	475.	500.	525.	550.	575.	.009	625.	650.	675.	700.	725.	750.	775.	800.	825.	850.	875.	006	925.	950.	975.

ABSORPTION COEFFICIENTS OF HZO (CONT.D) K (PER CM AT STP)

	3000 K	Ø	8E-	.109E-01	-	1	1	1	111	• 129E-01	.134E-01	l lil	44E-	50E-	57E-	64E-	72E-	BOE-	91E-	8	116-	L	33E-		10	67E-	76E-	87E-	.296E-01	05E-	.313E-01
	2500 K	0E-0	E-0	0E-0	.740E-02	40E-0	.760E-02	30E-0	20E-0	O-BOS	.910E-02	70E-0	03E-0	0	18E-0	24	32E-0	40	84	57	99	.174E-01	85	46	03	13	23	34	44	-	63
	2000 K	306-0	40E-0	50E-0	70E-0	0E-0	COE-0	20E-0		80E-0	L	70E-0	10E-0	60E-0	20E-0	80E-0	40E-0	0	0-360	17E-	24E-0	3	40E-	49	0	69E-	19	-306		12E-	N
	1500 X	0	50E-0	50E-0	S0E-0	0E-0	60E-0	70E-0	30E-0	90E-0	.210E-02	30E-0	60E-0	90E-0	30E-0	70E-0	10E-0	0	20E-0	90E-0	60E-0		10E-0	90E-0	70E-0	06E-0	16E-	N	35	4 SE-	
CM AT STP	000	00E-0	00E-0		OCE-O	.7COE-03	3	CUE-0	CE-0	COELO	10E-	0E-0	0E-0	0E-0	50E-0	0E-0	80E-3	00E-0	10E-0	40E-0	60E-0	L	10E-0	40E-0	70E-0	00E-0	0E-0	90E-0	50E-0	ō	0E-0
K (PER	0	60E-	4 SE-0	35	0E-0	65E-0	3	8E-0	7E-0	1		2 JE-0	S0E-0	85E-0	30E-0	80E-0	30E-0	95E-0	65E-0	50	45E-0	.840E-03	80E-0	OE-O	0E-0	0E-0	0E-0	90E-0	OE-	50E-0	80E-0
	3000 X	0-300	00E-0	00E-0	00E-0	00E-0	80E-0	30E-0	00E-0	00E-0	•185E-07	25E-0	60E-0	00E-0	80E-0	25E-0	70E-0	02E-0	80E-0	20E-0	80E-0	02E-0	80E-0	25E-0	70E-0	00E-0	70E-0	20E-0	00E-0	00E-0	21E-0
	1/CM	•000	025.	050	075.	100.	125.	150.	175.	200.	6225.0	250.	275.	300	325.	350.	375.	400	425.	450.	475.	500	525.	550.	575.	600	625.	<b>520</b>	675.	700.	725.

	3000 K	20	29		39		<b>1</b> 0	• 343E-01	.341E-01	<b>⊘</b>	33	.325E-01	13	.299E-01	83	63	43	23	90	66	00	87	0	4	4		63	9	.248E-01	0	• 202E-01
	2500 K	.272E-01	_	90E-	-386	56-	13E-	19ë-	22E-	2.7E-	0	30E-	4E-	14E-	03E-	91E-	79E-	63E-	47E-	28E-0	02	H	.202E-01	22E-	47E-	70E-	16	8E-	ທັ	27E-	•196E-01
(CONT.D)	2000 K	35E-	.247E-01	OE-	72E-	•285E-01	98E-		36-	34	39E-	4 2E	41	39	32E-	21	96E-	7	49E-	28E-	1 1	OOE-	22	48E-	73E-	.290E-01	82E-	58E-	32	0	•173E-01
NTS OF H20	1500 K	ហ	.174E-01		•195E-01	.208E-01	•224E-01	. 240E-01	62	74	05	25	4	58	60E-	55E-	38	1.	80	52	30	N	45	73	98	.310E-01	85	<b>V</b>	.202E-01	•164E-01	58
7 F	X 00,	III O	0 교	30E-0	를 당 1	·130E-01	ا لنا	67E-	<b>89</b> E-	128-		<b>65E</b>	.295E-01	40E	<b>306</b>	45E	이	ZOE	70E-	25E-	90	85E-	15	68E-	5E-	40E-	<b>60E</b> -	4E-	50	0E-	60E-0
BSO	! . ¥	30E-0	0-30	30E-0	0-300g	0	10E-0	BOE-0	0E-0	0	.128E-01	.152E-01	Φ	.222E-01	1	n	32E-	7	4	0	Ø	•475E-01	9	-	.418E-01	05-	5E-	0-35	.81 CE-02	0E-0	0
4	300	52E-0	85E-0	20E-0	555-0	90E-0	20E-0	60E-0	COE-0	60E-0	30E-0	208-0	60E-0	80E-0	32E-0	0ーヨンい	40E-0	88E-0	23E-0	70E-0	16E 0	26E 0	17E-0	40E-0	25E-0	40E-0	85E-0	82E-0	.170E-01	61E-0	45E-
	1/CM	750.	775.	800	825.	850.	875.	.006	925.	950.	975.	0000	025.	050	075.	100.	125.	150.	175.	200.	225.	250.	275.	300.	325.	350.	375.	400	7425.0	450.	475

		•	ABSORPTION	COEFFICIE	NTS OF H20	(CONT.D)		
1/CM	300	¥	O   X	1000 K	1500 K	2000 K	2500 K	3000 K
000	00E-	03	OE-O	40E-0	40E-0	40E-0	65E-	75E-0
25.	3000	00	OE O	0 300	00 300	12E-0	26E-	40E-0
30	OOE	00	OE C	OOE O	00E 0	70E-0	70E-0	07E-0
75.	OOE	00	OE O	OE O	OOE O	40E-0	00E-0	OE-
00	OOE	00	OE O	OOE O	OOE O	60E-0	25E-0	90E-0
25.	OOE	00	OE O	OOE O	OOE O	65E-0	37E-0	10E-0
50.	OOE	00	OE O	COEO	OOE O	20E-0	10E-0	60E-0
75.	OOE	00	OE O	OOE O	CEO	90E-0	60E-0	30
00	OOE	00	OE O	COEC	OOE O	70E-0	40E-0	10E-0
25.	OOE	00	OOE O	OOEO	OOE O	50E-0	22E-0	95E-0
50.	OOE	00	COEO	OOE O	OOE O	40E-0	15E-0	90E-0
75.	3000	0	OE O	OOE O	0 3000	30E-0	10E-0	90E-0
00	OOE	00	0 300	OOE O	OOE O	25E-0	0-360	90E-0
25.	OOE	0	OE O	0 300	OOE O	20E-0	03E-0	90E-0
50.	OOE	0	00E 0	COEC	OOE O	20E-0	05E-0	90E-0
75.	OOE	0	OE O	COEO	OOOE O	20E-0	06E-0	91E-0
000	OOE	00	OE O	COEO	OOE O	20E-0	07E-0	93E-0
20.00	OOE	00	OE O	00E 0	OCOE O	25E-0	115-0	97E-0
50	OOE	00	OE O	JOE O	OOOE O	E-0	(C)	.400E-02
75.	OOE	8	OE O	COEO	0000	34E-0	1 3E-0	01E-0
000	OOE	00	OE O	COEO	OOE O	40E-0	23E-0	0-390
25.	OOE	00	OE O	C 300	OOE O	46E-0	28E-0	10E-0
50	300	00	OE O	COE O	OCE O	51E-0	35E-0	19E-0
75.	OOE	0	0 3C	OOE O	00E 0	59E-0	44E-0	28E-0
00	OOE	00	OE C	OE O	0 300	65E-0	52E-0	38E-0
25.	OOE	00	OE C	COEO	OOE O	72E-0	58E-0	45E-0
50.	COOE	00	0 30	CCEO	OE O	81E-0	69E-0	57E-0
75.	OOE	00	OE C	COEC	OOE O	90E-0	19:-0	<b>ÚÚE−0</b>
8200.0	OOOE	00	• 000E 00	• OCCE 00	· 000E 00	S		·479E-02
25.	OOE	00	OE O	COEO	OE O	10E-0	98E-0	45E-0

		ABSORPTION K (PER	0 U	FFICI AT ST	ENTS OF	H20	(00 to 0)		
	Y O	O X	O	0	1 500	¥	2000 K	20 00 X	3000 K
	0	• 000E 0	0	OE C	• 000	0	20E-0	05E-0	90E-0
	E 00	.000E 00	8	0E 00	-000E	0	.330E-02	.415E-02	. 499E-02
w	0	• 000E 0	0	0 30	• 000	0	41E-0	21E-0	00E-0
u	0	0 3000°	0	0 30	000	0	55E-0	266-0	016-0
ш	0	• 000E 0	C	OE C	000	0	70E-0	376-0	026-
-	0	0 3000°	0	OE 0	000°	0	80E-0	40E-0	016-0
w	0	• 000E 0	0	0 30	000	0	88E-0	44E-0	00E-0
w	0	• 000E 0	00	OE O	• 000	0	91E-0	456-0	996-0
ш	0	• 000E 0	00	0 30	• 000	0	98E-0	46E-0	-366
Ä	0	• 000E 0	00	) 달	000	0	400E-0	45E-0	90E-0
~	0	0 3000°	0	0 30	000	0	00E-0	40E-0	80E-0
Ä	0	• 000E 0	00	OE O	• 000	0	3996-0	37E-0	756-0
~	0	• 000E 0	00	<b>○ 三</b> ○	000	0	396E-0	29E-0	635-0
ö	0	• 000E 0	000	0E 0	• 000	0	391E-0	23E-0	55E-0
ä	0	• 000E 0	၁	0E 0	000	0	391E-0	175-0	4 JR-0
ä	0	• 000E 0	၀	OE O	000	0	98E-0	166-0	いっぱーの
Ö	0	• 000E 0	00	이글이	• 000	0	400E-0	10E-0	<b>40E-0</b>
~	0	0 3000°	S	0 30	000	0	3996-0	04E-0	08E-0
<u>=</u>	0	o ≥000•	00	0 30	000	0	395E-0	95E-0	95E-0
ä	0	• 0000 o	$\circ$	OE O	000.	0	390E-0	85E-0	90E-0
2	0	O BOSO	00	0 30	• 000	0	3816-0	74E-0	67E-0
ä	0	0 3000°	0	O 글C	000	0	72E-0	62E-0	S1E-0
8	0	0 3000°	0	0	000	0	60E-0	48E-0	38E-0
2	0	0 3000°	0	E O	000	0	3486-0	36E-0	235-0
2	0	0 3000°	S	C L	000	0	316-0	20E-0	106-0
$\Xi$	0	C 3000.	0	0	000	0	136-0	04E-0	0- HG6
<u></u>	0	• 000E 0	O	0	• 000	0	94E-0	87E-0	80E-
7	0	•000E 0	O	m O	000.	0	73E-0	10E-0	• 268E-0
2	0	• 000E 0	0	O U	000.	0	516-0	344-0	こうだー
₩.	0	• 000E 0	O	S O	• • • •	0	30E-0	33E-0	40ki

	3000 X	2386-0	18E-0	072-0	197E-0	1876-0	1 79E-0	1716-0	1636	156E-0	149E-0	436-0	138E-0	1336-0	30E-0	1276-0	1266-0	1256-0	124E-0	1246-0	25E-0	126E-0	127E-0	1286-0	129E-0	306-0	1326-0	1326-0	1 25E-0	376-0	376-0
	2500 K	23E-0	01E-0	85E-0	1675-0	32E-0	142E-02	130E-02	.118E-02	09E-0	03E-0	80E-0	30E-0	20E-03	10E-03	00E-03	10E-03	10E-03	どうだーのふ	30F-03	40E-03	50E-03	BUE-03	00E-02	02E-02	04E-02	07E-02	106-02	136-02	16E-02	19E-0
(CONT.D)	2000 K	07E-0	83E-0	1625-0	40E-0	122E-0	04E-0	80E-0	.730E-03	620E-0	60E-0	530E-0	510E-0	500E-0	10E-0	530E-0	530E-0	360E-0	590E-0	610E-0	40E-0	660E-0	690E-0	20E-0	50E-0	80E-0	SOE-0	60E-0	00E-0	40E-0	80E-0
NTS OF H20	1500 K	OOOE O	OCOE O	OOOE O	OOOE O	OOOE O	OOCE O	OCOE O	• 000E 00	000E 0	000E 0	000E 0	OOOE O	000E 0	000E 0	OOOE O	DOOE O	OOOE O	COOE O	OOOE O	000E 0	000E 0	000E 0	OOOE O	OOOE O	000E 0	OOOE O	000E 0	OOOE O	SOOE O	200E 0
COEFFICIEN	1000 K	0000	COEO	OCOE O	0 300C	000E 0	0 3000	O BOCO	• 000E 00	OCOE O	000E 0	000E 0	000E 0	000E 0	0 3000	OOOE C	OCOE O	000E 0	000E 0	OOOE O	000E 0	000E 0	0 3000	000E 0	O BOOO	O BOOO	000E 0	000E 0	CCOE O	OOE O	OCOE O
ABSORPTION K (PER	У О	00E 0	OOE O	00E 0	OOE O	OOE O	OOE O	OOE O	• 000E 00	00E 0	00E 0	OOE O	OOEO	OOEO	OOE O	OOE O	OOEO	OOEO	OOE O	OOE O	OOE O	OOE O	OOEO	OOE O	OOE O	000	OCE O	OOEO	OOEO	00E 0	OOEO
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•	1/CM	•000	025.	080	075.	10c.	125.	150.	9175.0	200-	225.	250.	275.	300.	325.	350.	375.	400	425.	450.	475.	500	525.	550.	575.	600	625.	650.	675.	700.	725.

E 00 .000E 00 .102E-02 .12 E 00 .000E 00 .106E-02 .12 E 00 .000E 00 .113E-02 .13 E 00 .000E 00 .113E-02 .13
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(CONT.D)	2000 K	_ 0-1	·137E-02	E-0	16-0	7E-0	E-0		50E-0		80E-	40E-0	0E-0	20E-0	90E-3	80E-0	E-0	70E-0	60E-0	.260E-03	.250E-03	.250E-03
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